# Infinite Electrical Networks: Forward and Inverse Problems

Ian Zemke

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#### Abstract

In this paper we consider infinite electrical networks. Many of the results from finite electrical networks carry over in similar forms. In the infinite case, we show that the space of finite power voltages and current functions both form Hilbert spaces, which gives us powerful tools for analysis. Minimal power solutions come to the forefront of importance on these networks and in general we get minimal boundary-to-boundary maps. In a similar fashion to the finite case, we show how to recover a large class of infinite networks called "critical half-planar networks".

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# Chapter 1

# **Forward Problems**

# **1.1** Introduction

In this paper we develop the theory of infinite electrical networks. We first give some basic results about existence and uniqueness for boundary conditions in  $L^p$  for various p. It turns out that  $L^p$  spaces aren't the most natural spaces to work with for infinite graphs. Instead, we look at the space of functions of finite power, both for voltage functions and current functions. It turns out these are Hilbert spaces, which gives us some nice results. It turns out that certain subspaces of the finite power voltage functions and the finite power current functions are dual to each other in a natural sense. Perhaps most importantly, we prove various theorems relating to the existence of Dirchlet to Neumann maps and Neumann to Dirichlet maps and in what sense these maps exist for infinite graphs.

# **1.2** Preliminaries and Definitions

Firstly, we need some basic information about infinite graphs.

**Definition 1.2.1.** We call a graph **topologically connected** if there are no two subgraphs A and B such that  $A \cup B = G$  and  $A \cap B = \emptyset$ .

**Definition 1.2.2.** We call a graph **finitely connected** if every two vertices in the graph can be connected by a finite path.

Lemma 1.2.3. A graph G is topologically connected iff it is finitely connected.

*Proof.* If G is finitely connected then obviously it is topologically connected. Now suppose that G is topologically connected but not finitely connected and let v and w be vertices that can't be connected with a finite path. We will define a subgraph  $G_n(v)$ . A vertex a is in  $G_n(v)$  iff there is a path from v to a in G such that is of length n or less. An edge  $ab \in G_n$  iff  $a, b \in G_n$  and  $ab \in G$ . We will now define the graph

$$C(v) = \bigcup_{n \in \mathbb{N}} G_n(v)$$

which we will call the maximal component containing C(v). Now notice that  $\{C(v) : v \in G\}$  is just the collection of maximal finitely connected components of G, so if  $C(v) \neq C(w)$  (which occurs iff there is not a finite path from v to w) then  $C(v) \cap C(w) = \emptyset$ . Define the graphs

$$A = C(v)$$
 and  $B = \bigcup_{u \in \{G: C(u) \neq C(v)\}} C(u)$ 

and note that  $A \cap B = \emptyset$  but  $A \cup B = G$  so G is not topologically connected.  $\Box$ 

**Lemma 1.2.4.** If G is a finitely connected infinite graph with finite valence at each vertex, then G has countably many vertices and (at most) countably many edges.

*Proof.* Using the above notation, we have just seen that if  $v \in G$  is any vertex, then  $G = C(v) = \bigcup_{n \in \mathbb{N}} G_n(v)$ . Thus if we can show that each  $G_n(v)$  is finite we will be done. The rest of the proof will be by induction. Since every vertex has finite valence, if w is any vertex in G, we know that  $G_1(w)$  is finite. Now suppose that  $G_n(v)$  is finite for some  $n \geq 1$ . Then we have that

$$G_{n+1}(v) = \bigcup_{w \in G_n(v)} G_1(w)$$

which is a finite union of finite sets and is thus finite. By induction each  $G_n(v)$  is finite so the union is countable.

Now we need a definition of an infinite electrical network. We do this much the same as the finite case.

**Definition 1.2.5.** Let  $G = (\partial G, \operatorname{int} G, E)$  be an infinite graph divided into boundary and interior vertices  $\partial G$  and  $\operatorname{int} G$ . We will call the pair  $\Gamma = (G, \gamma)$  an infinite electrical network if the following conditions are satisfied

- 1. G is finitely connected,
- 2. each vertex of G has finite valence (i.e. finitely many edges connected to it),
- 3.  $\gamma$  is a function from E (the edges of the graph) to  $\mathbb{R}^+$  (strictly positive reals).

# **1.3** Naive Approaches to Infinite Graphs

In this section we develop some of the theory for  $L^p$  voltage functions, which is sort of the most naive route to approach this problem. We have some useful results, which we present for completeness, but the theory of  $L^p$  spaces in this context turns out to be not as robust as the theory of finite power voltages, which we will discuss later.

**Definition 1.3.1.** The spaces  $L^p(G)$  and  $L^p(E)$  are as follows

- 1. For  $1 \le p < \infty$   $L^p(G)$  is the space of real valued vertex functions u such that  $\sum_{v \in V(G)} |u(v)|^p < \infty$
- 2.  $L^{\infty}(G)$  is the space of bounded real valued vertex functions.
- 3.  $L^{p}(E)$  (resp.  $L^{\infty}(E)$ ) is the space of positive valued conductivity functions  $\gamma$  such that  $\sum_{e \in E} \gamma^{p} < \infty$  (resp. is bounded).

We recall that from Minkowski's inequality that these spaces are actually  $\mathbb{R}$ -vector spaces.

We note that since we assumed that vertices have finite valence, given an electrical network  $\Gamma = (G, \gamma)$  and a real valued function  $u : G \to \mathbb{R}$  and a, we can make sense of the current at each vertex. In fact, this lets us define will define the operator  $K : \mathbb{R}^G \to \mathbb{R}^G$  as

$$(Ku)(v) = \sum_{v' \sim v} (v - v')\gamma_{vv'}.$$

Thus a function u is called  $\gamma$  harmonic if (Ku)(v) = 0 for  $v \in \operatorname{int} G$ .

# **1.3.1** Results about $L^{\infty}(G)$

We can now state a result about existence.

**Theorem 1.3.2.** Let  $\Gamma$  be an infinite electrical network. Let  $\phi \in L^{\infty}(\partial G)$ . Then there exists a (not necessarily unique) function  $u \in L^{\infty}(G)$  such that  $u|_{\partial G} = \phi$  and u is  $\gamma$ -harmonic on int G. Furthermore, we have that  $||u||_{L^{\infty}(G)} = ||\phi||_{L^{\infty}(\partial G)}$ .

*Proof.* Arbitrarily pick a boundary vertex  $v_0$ . We can define the distance between two vertices v and v' as the minimum path length between v and v'. Let  $G_n$  denote the subgraph of G consisting of all vertices v with distance less than or equal to n from  $v_0$  and let an edge be between two vertices in  $G_n$  there is a corresponding edge in G. We can make  $G_n$  into an electrical network by setting  $\partial G_n = (\partial G \cap G_n) \cup (G_n \setminus G_{n-1})$  and defining a conductivity function on  $G_n$  to be just the restriction of  $\gamma$  onto  $G_n$ . We note that  $G_n$  is a finite graph. Define the function  $\phi_n : \partial G_n \to \mathbb{R}$  as  $\phi_n(v) = \phi(v)$  if  $v \in \partial G \cap G_n$  and set  $\phi_n(v) = 0$  if  $v \in \partial G_n \setminus \partial G$ . Using basic theory, the Dirichlet problem has a unique solution on finite graphs, so there is a function  $u_n : G_n \to \mathbb{R}$  such that  $u_n|_{\partial G_n} = \phi_n$  and  $u_n$  is  $\gamma$  harmonic on int  $G_n$ . We can extend  $u_n$  to all of G by setting u to be zero outside of  $G_n$ . We note that  $||u_n||_{L^{\infty}(G)} \leq ||\phi||_{L^{\infty}(G)}$  and hence  $\{u_n\}$  is a pointwise bounded sequence. Now order the vertices of G arbitrarily as the sequence  $\{v_i\}$ . Since  $u_n$  is bounded, we can find a subsequence  $u_{k_n^1}$  such that  $u_{k_n^1}(v_1)$  converges. Now find a subsequence  $\{u_{k_n^2}\}$  of this last sequence such that  $u_{k_n^2}(v_2)$  converges. Repeat this process to get a chain of subsequences such that  $u_{k_n^j}(v_j)$  converges. Now consider the diagonal subsequence  $u_{k_j^j}$ . This converges pointwise at every vertex, say to a function u. Now, since  $u_{k_j^j}$  is  $\gamma$ -harmonic on int  $G_\ell$  for all  $\ell \leq k_j$  and every interior vertex of G is contained in  $G_n$  for all n sufficiently large, we know that u will be  $\gamma$ -harmonic on G, and furthermore will obviously take the right boundary values. Finally, we note that since u has the right boundary values  $||u||_{L^{\infty}(G)} \geq ||\phi||_{L^{\infty}(\partial G)}$ . On the other hand, since  $||u_n||_{L^{\infty}(G)} \leq ||\phi||_{L^{\infty}(G)}$ , in the limit we will have  $||u||_{L^{\infty}(G)} \leq ||\phi||_{L^{\infty}(G)}$ .

### 1.3.2 Lack of Uniqueness

Unfortunately, uniqueness is in general not true, even if we only consider bounded functions. We might hope it would be true, since for instance in the half plane, harmonic functions that are zero on the real axis are zero on the upper half plane if we assume that they are bounded. The reader is advised to consider the case of an infinite string of conductors with a single conductor as a boundary vertex. By making the conductors have conductance  $2^n$  say, we can get  $\gamma$ -harmonic voltages of  $0, 1/2, 3/4, 7/8, \ldots$  which are bounded (the 0 voltage corresponds to the boundary vertex). Thus there's no hope of having uniqueness in this case without restrictions on  $\gamma$ . At the present time there are no known restrictions on  $\gamma$  to give uniqueness in the case of  $u \in L^{\infty}$ .

# **1.3.3** More on $L^p(G)$ spaces

It turns out that uniqueness is pretty easy easy in if we assume all of our functions are in  $L^p$  so that things go to zero as we move "far" into the graph, but the existence is much harder and we don't have any useful results. We will assume that the reader is somewhat familiar with measure spaces, and all measures will be assumed to be positive.

#### 1.3.4 Some Real Analysis

Here we develop a bit of machinery and terminology. We note that by Lemma 1.2.4, we can make G into a  $\sigma$ -finite measure space by simply putting the counting measure with weight 1 on each vertex. The set of measurable sets is just  $\mathcal{P}(V(G))$ , the set of all subsets of V(G) (the vertices of G). We begin with some useful but extremely basic remarks.

**Lemma 1.3.3.** Let  $(X, M, \mu)$  be a measure space and let  $E_1 \subseteq E_2 \subseteq \cdots$  be an increasing sequence of measurable sets. If  $f \in L^p(X)$ , then  $\int_{E_i \setminus E_{i-1}} |f|^p \to 0$ .

*Proof.* This is super trivial. Define  $F_1 = E_1$  and let  $F_n = E_n \setminus E_{n-1}$  for  $n \ge 2$ . Note that the collection  $\{F_n\}$  is a pairwise disjoint collection of sets and hence, by the definition of a measure we have that

$$\int_E |f|^p = \sum_j \int_{F_j} |f|^p \le \int |f|^p < \infty$$

and since each  $\int_{F_j} |f|^p$  is nonnegative, we have absolute convergence and hence the summands must tend to zero as  $j \to \infty$ .

**Lemma 1.3.4.** In the case of G, which we treat as a measure space with the counting measure, we have that  $||f||_{L^{\infty}} \leq ||f||_{L^{p}}$ .

This is obvious since the weights on each vertex are 1.

# **1.3.5** First Results for $L^p(G)$

**Theorem 1.3.5** (Maximum Principle for Infinite Graphs). If u is  $\gamma$ -harmonic on  $\Gamma$  and  $u \in L^p(G)$ , then  $\|u\|_{L^{\infty}(G)} = \|u|_{\partial G}\|_{L^{\infty}(\partial G)}$ .

*Proof.* This is basically just an application of the maximum principle for finite graphs. Let  $u \in L^p(G)$  for  $1 \leq p < \infty$  and let  $\phi$  denote  $u|_{\partial G}$ . Pick a vertex  $v \in G$  arbitrarily and let  $G_n = G_n(v)$  be the finite electrical network as defined in our discussion of  $L^{\infty}$ . Let  $\partial G_n$  and int  $G_n$  also be as defined above. Note that  $G_1 \subseteq G_2 \subseteq \cdots$  is an increasing sequence of sets. Hence by Lemma 1.3.3 we know that

$$\int_{G_{n+1}\backslash G_n} |u|^p \to 0.$$

By Lemma 1.3.4 this implies that  $||u||_{L^{\infty}(G_{n+1}\setminus G_n)} \to 0$ . Since u is  $\gamma$ -harmonic on  $G_n$ , by the maximum principle for finite graphs, we know that

$$\|u\|_{L^{\infty}(G_n)} \le \|u\|_{L^{\infty}(\partial G_n)}.$$
(1.1)

We recall that by the definition of  $\partial G_n$ , we have that

$$\partial G_{n+1} = (\partial G \cap G_{n+1}) \cup (G_{n+1} \setminus G_n) \subseteq \partial G \cup (G_{n+1} \setminus G_n).$$

Hence

$$\|u\|_{L^{\infty}(\partial G_n)} \le \|u\|_{L^{\infty}(\partial G)} + \|u\|_{L^{\infty}(G_{n+1}\setminus G_n)}.$$
(1.2)

Clearly  $||u||_{L^{\infty}(G_n)} \to ||u||_{L^{\infty}(G)}$ . Combining equations (1.1) and (1.2) we get that

$$||u||_{L^{\infty}(G_n)} \leq ||u||_{L^{\infty}(\partial G)} + ||u||_{L^{\infty}(G_{n+1}\setminus G_n)}.$$

Letting  $n \to \infty$  and using the various results about convergence of various terms shows that

$$\|u\|_{L^{\infty}(G)} \le \|u\|_{L^{\infty}(\partial G)}$$

Since the reverse inequality is trivial, we have equality in the above expression.  $\hfill\square$ 

**Theorem 1.3.6.** Let  $u_1, u_2 \in L^p(\partial G)$  for  $1 \leq p < \infty$  and  $u_1|_{\partial G} = u_2|_{\partial G}$ . Then  $u_1 = u_2$ .

*Proof.* Consider the function  $h = u_1 - u_2$ . Note that h is  $\gamma$ -harmonic,  $h \in L^2$  and h = 0 on  $\partial G$ . By the maximum principle for infinite graphs, we know that  $\|h\|_{L^{\infty}(G)} \leq \|h\|_{L^{\infty}(\partial G)} = 0$  and hence h = 0 on G so  $u_1 = u_2$ .

### **1.3.6** Existence Theorems for $L^p(G)$

In general, we do not have existence in  $L^p(G)$ , for example take an infinite series of conductors with unit conductivity on each edge and a single boundary vertex (the exact configuration of these conductors is not that important, they can extend in both directions from the boundary vertex or they can extend in only one direction). Set voltage 1 on the boundary. Clearly this is in  $L^p(G)$  for all p. But for  $p \neq \infty$ , we readily see that no  $L^p$  solution will exist, since every nonconstant  $\gamma$ -harmonic function will be unbounded while the constant voltage  $\phi = 1$  will not be in  $L^p$ .

# **1.4** The Spaces of Finite Power Functions

It turns out the most natural condition to consider for functions on an infinite graph is the space of voltages that satisfy finite power. This turns out to be much more natural than  $L^p$  spaces since it takes into account the conductivities better than it seems is possible for  $L^p$  spaces. It turns out that there are sort of two dual finite power spaces. There is the space of finite power voltage functions, and there is the space of finite power current functions, which are basically be functionally dual to each other.

#### 1.4.1 Finite Power Voltage Functions

**Definition 1.4.1.** If  $\Gamma$  is an infinite resistor network and  $\phi$  is a real valued vertex function, we define the power of  $\phi$  to be

$$P(\phi) \stackrel{\text{def}}{=} \sum_{v \in G} \sum_{v' \sim v} \gamma_{vv'} (\phi(v) - \phi(v'))^2.$$

We note that by convention  $\gamma_{vv'} = 0$  iff  $v \not\sim v'$  and hence we can write the above sum as

$$\sum_{(v,v')\in V\times V} \gamma_{vv'}(\phi(v) - \phi(v'))^2$$

or just

$$\sum_{V \times V} \gamma_{vv'} (\phi(v) - \phi(v'))^2$$

for brevity.

We note that there is no reason to assume that  $P(\phi)$  is finite for a particular  $\phi$  or even a  $\gamma$  harmonic  $\phi$ .

**Definition 1.4.2.** We define  $F(\Gamma)$  to be the set of real valued vertex functions on  $\Gamma$  of finite power.

**Lemma 1.4.3.**  $F(\Gamma)$  is a vector space over  $\mathbb{R}$ .

*Proof.* Clearly  $F(\Gamma)$  is closed under multiplication by scalars. That  $F(\Gamma)$  is closed under addition is just the triangle inequality for  $L^2$  since

$$\begin{split} \sqrt{P(f+g)} &= \left(\sum_{V \times V} \gamma_{vv'}(f(v) + g(v) - f(v') - g(v')^2\right)^{1/2} \\ &= \left(\sum_{V \times V} \left(\sqrt{\gamma_{vv'}}(f(v) - f(v')) + \sqrt{\gamma_{vv'}}(g(v) - g(v')\right)^2\right)^{1/2} \\ &\leq \left(\sum_{V \times V} (\sqrt{\gamma_{vv'}}(f(v) - f(v'))^2\right)^{1/2} + \left(\sum_{V \times V} (\sqrt{\gamma_{vv'}}(g(v) - g(v'))^2\right)^{1/2} \\ &= \sqrt{P(f)} + \sqrt{P(g)}. \end{split}$$

It turns out that we can in some sense solve the Dirichlet problem in  $F(\Gamma)$ , but that  $F(\Gamma)$  isn't quite the right space to look at. We note that adding a constant to every vertex does not change the power, and that the constant function has 0 power, this leads us to make the following definition:

**Definition 1.4.4.** We define the space  $Z(\Gamma)$  to be  $F(\Gamma)/\text{span}\{1\}$  where 1 denotes the constant function 1.

We note that the power function is well defined on  $Z(\Gamma)$ . We recall that power was almost a norm on  $F(\Gamma)$  but on  $Z(\Gamma)$  it turns out to be a norm:

**Theorem 1.4.5.**  $Z(\Gamma)$  is a Hilbert space with inner product given by

$$(f,g) = \sum_{v,v'} \gamma_{vv'}(f(v) - f(v'))(g(v) - g(v')).$$

Proof. The only claim that requires justification is that  $Z(\Gamma)$  with the inner product above is Cauchy complete. We use the standard theorem that says a normed vector space is complete iff  $\sum_{\mathbb{N}} ||f_n|| < \infty$  implies  $\sum_{\mathbb{N}} f_n$  exists as a limit in the norm topology (Theorem 5.1 of Folland). Thus suppose that  $\sum_{\mathbb{N}} \sqrt{P(f_n)} < \infty$ . We first claim that we get pointwise convergence in a certain sense. Let  $v_0$  be an arbitrary vertex in G and pick representatives of  $f_n$  in  $F(\Gamma)$ such that  $f_n(v_0) = 0$ . We first claim that  $\sum_n |f_n(v)| < \infty$  for all  $v \in G$ . This part of the proof will be by induction. Suppose the claim holds for all vertices of distance k or less from  $v_0$  and consider a vertex  $v_{k+1}$  of distance exactly k+1. Let  $v_0v_1 \dots v_kv_{k+1}$  be a path from  $v_0$  to  $v_{k+1}$ . So by assumption

$$\sum_{n\in\mathbb{N}} |f_n(v_{k+1})| < \infty.$$
(1.3)

Since  $\sum_{n \in \mathbb{N}} \sqrt{P(f_n)} < \infty$ , we in particular have that

$$\sqrt{\gamma_{v_k v_{k+1}}} |f_n(v_k) - f_n(v_{k+1})| \le \sqrt{P(f_n)}$$

and hence

$$\sum_{n \in \mathbb{N}} |f_n(v_k) - f_n(v_{k+1})| < \infty.$$

Applying the triangle inequality shows that

$$\sum_{n=1}^{N} |f_n(v_{k+1})| - \sum_{n=1}^{\infty} |f_n(v_k)| \le \sum_{n=1}^{\infty} |f_n(v_k) - f_n(v_{k+1})|.$$

Letting  $N \to \infty$  and using equation (1.3), we see that  $\sum_{n \in \mathbb{N}} |f_n(v_{k+1})| < \infty$ . By induction on k we thus know that  $\sum_n f_n(v)$  converges absolutely for each v. Denote the limiting function by f. Firstly, we note that  $f \in F(\Gamma)$ , since an application of Fatou's lemma shows that

$$\sqrt{P(f)} = \left(\sum_{V \times V} \gamma_{vv'} (f(v) - f(v'))^2\right)^{1/2}$$
$$\leq \lim_{n \to \infty} \left(P\left(\sum_{i=1}^n f_i\right)\right)^{1/2} \leq \sum_{i=1}^\infty \sqrt{P(f_i)}.$$

To see that  $\sum_{n=1}^{N} f_n$  converges to f in the power norm, we will apply the dominated convergence theorem. Firstly, let  $G_N(v, v') = \sum_{i=1}^{N} |f_n(v) - f_n(v')|$  and let  $G(v, v') = \sum_{i=1}^{\infty} |f_n(v) - f_n(v')|$ . We note that G is finite everywhere by the triangle inequality and the fact that  $\sum_{n \in \mathbb{N}} |f_n(v)| < \infty$  for all  $v \in G$ . We first claim that  $\sum_{v,v'} \gamma_{vv'} (G(v, v'))^2$  is finite. To see this, note that  $G_n$  is pointwise nondecreasing in n and

$$\left(\sum_{V \times V} \gamma_{vv'} \left(\sum_{i=1}^{N} |f_n(v) - f_n(v')|\right)^2\right)^{1/2} \le \sum_{i=1}^{N} \left(\sum_{V \times V} \gamma_{vv'} (f_n(v) - f_n(v')^2\right)^{1/2}$$

by the triangle inequality. By assumption the latter sum is finite (since it is just  $\sum_{\mathbb{N}} \sqrt{P(f_n)}$ ) and hence by the Monotone Convergence Theorem, we know that

$$\sum_{v,v'} \gamma_{vv'} (G(v,v'))^2 = \lim_{n \to \infty} \sum_{v,v'} \gamma_{vv'} (G_n(v,v'))^2 < \infty.$$
(1.4)

We use the estimate  $(a + b)^2 \leq 4(a^2 + b^2)$  to get that

$$\gamma_{vv'} \left| \left( \sum_{i=1}^{N} f_n(v) - f_n(v') \right) - \sum_{i=1}^{\infty} (f_n(v) - f_n(v')) \right|^2$$

$$\leq 4\gamma_{vv'} \left| \sum_{i=1}^{N} (f_n(v) - f_n(v')) \right|^2 + \left| \sum_{i=1}^{\infty} (f_n(v) - f_n(v')) \right|^2$$
$$\leq 4\gamma_{vv'} \left( \sum_{i=1}^{N} |f_n(v) - f_n(v')| \right)^2 + \left( \sum_{i=1}^{\infty} |f_n(v) - f_n(v')| \right)^2$$
$$\leq 8\gamma_{vv'} G(v, v')^2$$

which is in  $L^1(V \times V)$  by the estimate on in (1.4). Thus by the dominated convergence theorem, we know that

$$P\left(\sum_{i=1}^{N} f_{i} - \sum_{i=1}^{\infty} f\right) = \sum_{V \times V} \gamma_{vv'} \left| \left(\sum_{i=1}^{N} f_{n}(v) - f_{n}(v')\right) - \sum_{i=1}^{\infty} (f_{n}(v) - f_{n}(v')) \right|^{2} \to 0$$

since the integrand goes to zero pointwise and is pointwise bounded by  $8\gamma_{vv'}G(v,v')^2 \in L^1(V \times V)$ . Hence  $Z(\Gamma)$  is Cauchy complete.

### 1.4.2 Interlude about Hilbert Spaces

We now need to present some basic machinery about Hilbert spaces in order to discuss the space of finite power functions. Many of the existence and uniqueness theorems that we will state and prove can be proven (and indeed were initially proven) without using the language of Hilbert spaces, but it turns out that by introducing some basic machinery, we can simplify many of the proofs significantly. We begin with a basic result, the proof of which is left as a reference.

**Lemma 1.4.6** (from pg 175 of [3]). If M is a closed subspace of a Hilbert space H then  $H = M \oplus M^{\perp}$ , that is each  $x \in H$  can be uniquely expressed as x = y+z where  $y \in M$  and  $z \in M^{\perp}$ . Moreover, y and z are the unique elements of M and  $M^{\perp}$  whose distance to x is minimal.

See Folland's book on Real Analysis for the proof, though the theorem is extremely standard in the subject of Hilbert spaces

**Lemma 1.4.7.** If *H* is a Hilbert space and *M* is a closed affine subspace, then there is a unique element *x* of *M* such that ||x|| is minimal in *M*.

*Proof.* The proof follows essentially from the last lemma, in fact only from the last part of the last lemma. By an affine subspace, we mean that M-y is a linear subspace for some  $y \in H$ . By the Lemma, there is a unique  $x \in M - y$  such that the distance from x to -y is minimal in M-y. Hence ||x + y|| is minimal over M - y. But x = z - y where  $z \in M$ , but ||x + y|| = ||z|| so clearly ||z|| is minimal over M, as we wanted. Also z is clearly the unique minimizer.

#### 1.4.3 More on Finite Power Voltage Functions

We are now in a position to state several theorems about existence and uniqueness of the Dirichlet problem.

**Definition 1.4.8.** We will denote the set of functions in  $Z(\Gamma)$  which are constant on  $\partial G$  by W.

**Lemma 1.4.9.** The set W is a closed subspace of  $Z(\Gamma)$ .

*Proof.* Obviously W is a subspace, so it is only necessary to show that W is closed. We note that convergence of a sequence in  $Z(\Gamma)$  implies pointwise convergence of that sequence (in the sense that we can pick representatives of the sequence which converge pointwise) as was shown in the proof that  $Z(\Gamma)$  was Cauchy complete. Hence the property of being constant on the boundary will be preserved under limits. Hence W is a closed subspace.

**Lemma 1.4.10.** Let  $\phi \in Z(\Gamma)$ . Then there is a function  $u \in Z(\Gamma)$  such that P(u) is minimum over all functions  $f \in Z(\Gamma)$  such that  $f|_{\partial G} = \phi|_{\partial G}$ .

*Proof.* Since  $\phi + W$  is a closed affine subspace, we apply Lemma 1.4.7 to see that there is a unique element u of  $\phi + W$  of minimal norm and furthermore, we have that  $u|_{\partial G} = \phi|_{\partial G}$ . Lastly we note that every function  $f \in Z(\Gamma)$  that agrees with  $\phi$  on the boundary will be in  $\phi + W$ , so the statement of the theorem follows.

We now show that these minimal solutions are  $\gamma$ -harmonic.

**Lemma 1.4.11.** Suppose that  $\phi \in Z(\Gamma)$  and that u is of minimal norm in  $\phi + W$ . Then u is  $\gamma$ -harmonic on int G.

*Proof.* If  $\phi$  is of minimal norm in  $\phi+W$  then by basic facts about Hilbert Spaces, since W is a closed subspace of  $Z(\Gamma)$  we have that  $\phi \in W^{\perp}$ . In particular, the set indicator function  $\chi_v$  is in W for all  $v \in \operatorname{int} G$ . Hence  $(\phi, \chi_v) = 0$ , but a trivial computation shows that

$$(\phi, \chi_v) = 2 \sum_{v' \sim v} \gamma_{vv'}(\phi(v) - \phi(v')),$$

and since this is zero for all  $v \in \operatorname{int} G$  we know that  $\phi$  is  $\gamma$ -harmonic.

**Theorem 1.4.12.** If  $\phi \in Z(\Gamma)$ , then there is a unique element  $u \in Z(\Gamma)$  such that u is minimal in power in  $\phi + W$ . Furthermore, u is  $\gamma$ -harmonic on int G.

This is literally just a restatement of the previous two theorems. We note that this is in some ways an existence and uniqueness result for the Dirichlet problem. But we should note that it is not a solution as we may want. This gives us a map from valid boundary voltages to a currents coming from a minimal power solution, but in general, there may be other finite power  $\gamma$ -harmonic functions with the same boundary voltages but different currents leaving the boundary.

**Remark 1.4.13.** The Dirichlet problem as it is often formulated in the finite case is ill posed, even in  $Z(\Gamma)$ . In particular, functions in  $Z(\Gamma)$  with the same boundary voltages may have different current boundary currents.

We should note that we only have a unique energy minimizing solution and not a unique  $\gamma$ -harmonic solution. In fact, even the assumption of finite power does not imply that there is a unique finite power  $\gamma$ -harmonic function with given boundary voltages. An example is an infinite string of conductors in series, with conductance such that  $\sum_{\mathbb{Z}} 1/\gamma^2 < \infty$ , and a single boundary vertex in the middle.



Figure 1.1: The infinite series of conductors discussed in Remark 1.4.13.

Consider the voltage function with constant current 1 flowing constantly the right on the entire graph. This will give us a finite power voltage function on the entire graph since the power can be rewritten as  $\sum I/\gamma^2 = 1 \sum 1/\gamma^2$ . There is no current flowing out of the network with this voltage. Note that we can alter the above example to have constant 1 current flowing to the right, on the left side of the graph, and then have constant 1 current leaving the boundary vertex, and no current flowing on the right side of the graph. These are two  $\gamma$ -harmonic finite power solutions with different boundary currents. This fact will turn out to be important later: there is not in general a Dirichlet map from boundary voltages to boundary currents.

#### 1.4.4 Characterizing Minimal Power Solutions

We continue with our discussion of the the minimal power solutions of  $Z(\Gamma)$ . We note that if  $\phi \in Z(\Gamma)$ , by the above theorem, we can find a  $u \in Z(\Gamma)$  such that u has minimal energy and  $u|_{\partial G} = \phi|_{\partial G}$ . Thus given a  $u \in Z(\Gamma)$  we can determine whether it has the property of having minimal power with respect to its boundary conditions, i.e. whether u is the element of least norm in the space u + W.

**Definition 1.4.14.** The set of  $u \in Z(\Gamma)$  which have minimal power will be denoted by  $M(\Gamma)$ .

**Lemma 1.4.15.**  $M(\Gamma)$  is a subspace of  $Z(\Gamma)$ .

*Proof.* The proof is remarkably straightforward, and extremely enlightening. Obviously scaling by constants doesn't change the property of being minimal. We note that the minimal power functions correspond exactly to minimal norm elements in closed affine subspaces of  $Z(\Gamma)$ . To see this, note that if  $f \in M(\Gamma)$ then f has minimal norm over all elements u in  $Z(\Gamma)$  such that f-u is supported only in the interior of G (or to be more exact is constant on  $\partial G$ ). Recall that the space of functions that are constant on the boundary is denoted by W. Then f is minimal iff f has minimal norm in f + W, but this happens iff  $f \in W^{\perp}$ . We note that  $f, g \in M(\Gamma)$  iff  $f, g \in W^{\perp}$  which implies that  $f + g \in W^{\perp}$  which occurs iff  $f + g \in M(\Gamma)$  so  $M(\Gamma)$  is a subspace.

As a corollary of the proof of the previous theorem, we have that

**Corollary 1.4.16.** If W denotes the subspace of  $Z(\Gamma)$  consisting of elements which are constant on  $\partial G$ , then  $M(\Gamma) = W^{\perp}$ . Since W is closed we know that  $W^{\perp} = M(G)$  is a closed subspace.

### 1.4.5 Finite Power Current Functions

We will discuss the Neumann problem in much the same way that we discussed the Dirichlet problem above. We will define a space of finite power current functions, which will turn out to be a Hilbert space,

**Definition 1.4.17.** Let  $\phi \in F(\Gamma)$ . Then we define the function  $I(\phi)$  to be the real valued directed edge function defined by

$$I(\phi)(vv') = \gamma_{vv'}(\phi(v) - \phi(v')).$$

**Definition 1.4.18.** We define a current function I on the directed edges of a graph to be a function such that I(vv') = -I(v'v) and such that  $\sum_{v'\sim v} I(vv') = 0$  for all  $v \in \operatorname{int} G$ .

We note that  $I(\phi)$  is a current function iff  $\phi$  is  $\gamma$ -harmonic on int G.

**Definition 1.4.19.** We recall the definition of the map  $K : Z(\Gamma) \to \mathbb{R}^G$ . We define analogously the map  $\widetilde{K} : F(E) \to \mathbb{R}^G$  define by

$$\widetilde{K}(I)(v) = \sum_{v' \sim v} I(vv').$$

**Definition 1.4.20.** We define the power of a current function I to be

$$\sum_{vv':\gamma_{vv'}\neq 0} \frac{I(vv')^2}{\gamma_{vv'}}$$

**Definition 1.4.21.** We define the space of finite power current F(E) to the be set of current functions with finite power.

**Lemma 1.4.22.** F(E) is a Hilbert space with inner product

$$(I_1, I_2) = \sum_{v, v': \gamma_{v, v'} \neq 0} \frac{I_1(vv')I_2(vv')}{\gamma_{v, v'}}$$

Proof. As with the case of voltages, the only point worth mentioning is Cauchy completeness. The proof of this fact is similar to the proof of Cauchy completeness of  $L^2$  for a general measure space. We observe that if we let d# be the counting measure on  $V \times V$ , then the above inner product is exactly the inner product on  $L^2(V \times V, \mu)$  where  $d\mu = \gamma(v, v')d\#$ . Hence to show that F(E) is a Cauchy complete, it is sufficient to show that F(E) is closed in  $L^2(V \times V, \mu)$ . If  $I_n$  is Cauchy, it will converge to a function  $\tilde{I} \in L^2(V \times V, \mu)$ . To show that the limiting function is a current function, we note that in general, convergence in  $L^p$  for  $1 \leq p < \infty$  for any counting measure implies pointwise convergence on sets of positive measure. Hence we will have that  $I_n \to \tilde{I}$  pointwise, which will imply that  $\tilde{I}(vv') = \tilde{I}(v'v)$  and that the sum of currents coming into a vertex will be zero since it is zero for each  $I_n$ . Hence F(E) is a closed subset of a Cauchy complete space and is hence Cauchy complete.

**Theorem 1.4.23.** Let  $i_0 \in F(E)$ . Then there is a unique current function i such that P(i) is minimal over all functions  $i' \in F(E)$  such that  $\widetilde{K}i' = \widetilde{K}i_0$ . Furthermore, there is a unique  $u \in Z(\Gamma)$  such that I(u) = i.

*Proof.* Let Y denote the space of finite power current functions i such that  $\tilde{K}i = 0$ . We note that i is a closed subspace since convergence in F(E) implies pointwise convergence since the topology on F(E) is the same as the subspace topology given from the topology on  $L^2(V \times V)$  under the appropriate counting measure and convergence in  $L^p$  under a counting measure (for  $1 \leq p < \infty$ ) implies pointwise convergence since each point has positive measure. But as in the case of  $Z(\Gamma)$ , by 1.4.7 we know that there is always a unique element of  $i_0 + Y$  of minimal norm. Call this element i.

Now we just need to establish that there is a unique  $u \in Z(\Gamma)$  such that I(u) = I. Uniqueness is fairly straightforward, since if  $u_1$  and  $u_2$  both satisfy  $I(u_1) = I(u_2)$  then  $I(u_1-u_2) = 0$  and hence there is no current flowing anywhere for the voltage function  $u_1 - u_2$  and hence  $u_1 - u_2$  is constant and hence equal to zero in  $Z(\Gamma)$ . The existence of such a u is harder, but not overly difficult. It relies on an argument made by Will Johnson. It is sufficient to show that if the sum of  $I(ij)^2/\gamma_{ij}$  around any loop in G is zero, then such a u will exist. This is clear since if the sum around every loop is zero, then we can just define u by arbitrarily setting u to be zero at a single vertex, and then defining u by extending u to neighboring vertices by defining it to so that the voltage difference yields the desired current. This is well defined if the sum around a loop of the necessary voltage differences is zero.

Let  $C = v_0 v_1 v_2 \cdots v_n v_0$  be a loop in G. Let  $i_C$  have current have current one along each edge  $v_k v_{k+1}$  and have -1 along each edge  $v_k v_{k-1}$  (where say  $v_{-1} = v_n$  and  $v_{n+1} = v_0$ ). We note that  $\widetilde{K}i_C = 0$  and hence  $i_C \in Y$ . But by 1.4.7 we know that  $i \in Y^{\perp}$  and hence  $(i, i_C) = 0$ . But, we just observe by explicit computation that

$$(i, i_C) = 2\sum_{j=0}^{n} \frac{i(v_i v_{j+1})}{\gamma_{v_j v_{j+1}}}$$

which must be zero, exactly as we needed. Hence there is a  $\phi$  such that  $I(\phi) = i$ . We should also note that clearly  $\phi \in Z(\Gamma)$  since  $P(\phi) = P(i)$  by definition.  $\Box$ 

This formulation of the Neumann problem is somewhat different than in the finite graph case. We should note that something happens which is unexpected:

**Remark 1.4.24.** There is not a well defined Neumann-to-Dirichlet map in general, even if we restrict ourselves to functions in  $Z(\Gamma)$ .

In particular, given boundary currents on a network G, even if we know that they came from a voltage function in  $Z(\Gamma)$ , we cannot say that they came from a unique function. For example, take a countably infinitely many conductors in series as in Figure 1.1 and suppose that the conductivities satisfy  $\sum_i \frac{1}{\gamma} < \infty$ (for instance take  $\gamma_n = n^2$ ). Take a vertex and arbitrarily set the voltage u to zero. Assume that the current is constant on the entire graph, say 1 to the right. Clearly with this information we can extend  $u \gamma$ -harmonically to the rest of the graph, but we see that the power on the graph is  $\sum_{\mathbb{Z}} \frac{I^2}{\gamma} = 1^2 \sum_{\gamma} \frac{1}{\gamma} < \infty$ ! If we set any subset of the graph to be boundary vertices, we note that the current leaving the network is 0 at each boundary vertex, even though the solution is not constant.

#### 1.4.6 Minimal Currents

Just like in the case of minimal power voltage functions in  $Z(\Gamma)$ , we can consider minimal current functions, and we will analogous results.

**Definition 1.4.25.** Let M(E) denote the set of minimal current functions in F(E).

**Lemma 1.4.26.** M(E) is a closed subspace of F(E). Moreover, if Y denotes the set of all current functions i in F(E) with  $\tilde{K}i = 0$ , then M(E) is exactly  $Y^{\perp}$ .

*Proof.* As in the case of voltage functions, we note that i is minimal iff i has minimal norm over the set of current functions i' such that  $\widetilde{K}i' = \widetilde{K}I$ . The set of such functions is exactly i + Y. Hence i has minimal norm over i + Y iff  $i \in Y^{\perp}$ .

#### 1.4.7 Duality of Current and Voltages

We now have two interesting spaces, namely the space of minimal voltage functions  $M(\Gamma)$ , and the space of minimal current functions M(E). Since all countably infinite dimensional Hilbert spaces are isometrically isomorphic, it is too easy to say that  $M(\Gamma) \cong M(E)$  since that is trivially true. They turn out to not be naturally isomorphic under the obvious maps, so instead, we try to see how close these spaces are isomorphic in the *natural* way. We need some definitions. **Definition 1.4.27.** Let  $H(\Gamma)$  denote the subspace of  $Z(\Gamma)$  of functions which are  $\gamma$ -harmonic on int G. Similarly let L(E) denote the subspace of F(E) consisting of current functions such that the sum of the necessary potential differences in any loop is zero.

We note that  $M(\Gamma) \subseteq H(\Gamma)$  and  $M(E) \subseteq L(E)$ . We have natural functions between L(E) and  $H(\Gamma)$ .

**Definition 1.4.28.** Let  $I : H(\Gamma) \to L(E)$  denote the function which takes a  $\gamma$ -harmonic function in  $Z(\Gamma)$  and maps it to the natural current function induced. Similarly let  $\Phi : L(E) \to H(\Gamma)$  denote the function which takes a current function with necessary loop sums that are zero and maps it to the unique voltage function that satisfies those currents.

**Lemma 1.4.29.** Both *I* and  $\Phi$  are bijective linear functions. Furthermore both are isometries. Lastly  $I = \Phi^{-1}$ .

*Proof.* The fact that they are isometries is trivial. To see that I and  $\Phi$  are bijective, we just observe that given a function in  $I \in L(E)$ , we have already shown that there is a unique function  $u \in Z(\Gamma)$  such that I(u) = I, and hence we clearly have that  $I\Phi = id|_{L(E)}$  and  $\Phi I = id|_{H(\Gamma)}$  so the claim follows.  $\Box$ 

Hence we have the following diagram illustrating the relationship between the respective sets concerning voltages and currents:

$$\begin{array}{ccc} Z(\Gamma) & F(E) \\ \cup & \cup \\ H(\Gamma) & \stackrel{I = \Phi^{-1}}{\longleftrightarrow} & L(E) \\ \cup & \cup \\ M(\Gamma) & M(E) \end{array}$$

A natural conjecture would be that I and  $\Phi$  map minimal elements to minimal elements. This is unfortunately not true. A counterexample to  $\Phi$ mapping minimal currents to minimal voltages is again offered by an infinite series of conductors with a single boundary vertex in the middle. There will be a minimal current such that the current leaving the system is nonzero, while any minimal voltage will just be constant, and hence there are strictly more minimal currents in this example than minimal voltages. We will later see that I cannot in general map minimal voltages to minimal currents.

We should note that in general, we can imagine evaluating I at voltages which are not in  $H(\Gamma)$ , but it wouldn't immediately be obvious how to evaluate  $\Phi$  on edge functions which are not in L(E). We will use Hilbert space duality to get around this problem. We remark that there is a sort of symmetry between current and voltage, which leads us to define the following function.

**Definition 1.4.30.** We will define the pseudoinner product  $G : Z(\Gamma) \times F(E) \to \mathbb{R}$  defined by

$$G(u,i) = \sum_{V \times V} (u(v) - u(v'))i(vv').$$

We note that G(u, i) is always finite since by Cauchy Schwarz we know that

$$\sum_{V \times V} |(u(v) - u(v'))i(vv')| = \sum_{V \times V} \sqrt{\gamma_{vv'}} |(u(v) - u(v'))| \frac{|i(vv')|}{\sqrt{\gamma_{vv'}}}$$
$$\leq \left(\sum_{V \times V} \gamma_{vv'}(u(v) - u(v'))^2\right)^{1/2} \left(\sum_{V \times V} \frac{|i(vv')|^2}{\gamma_{vv'}}\right)^{1/2}$$

which the understanding that the integrand is 0 whenever  $v' \not\sim v$ . The above computation shows immediately that

$$|G(u,i)| \le ||u||_{Z(\Gamma)} ||i||_{L(E)}.$$
(1.5)

We now recall the Riesz Representation theorem:

**Theorem 1.4.31.** Let H be a Hilbert space and let  $f : H \to \mathbb{R}$  be a bounded linear functional. Then there is a unique  $y_f \in H$  such that  $f(x) = (y_f, x)$  for all  $x \in H$ . Furthermore,  $||f||_{\sup} = ||y_f||$ .

Thus we have the following result:

**Lemma 1.4.32.** The function G defined above yields well defined, bounded linear functions  $\tilde{I}$  and  $\tilde{\Phi}$  defined on  $Z(\Gamma)$  and F(E) respectively such that  $\tilde{I}|_{H(\Gamma)} = I$  and similarly  $\tilde{\Phi}_{L(E)} = \Phi$ .

Proof. For a fixed  $u \in Z(\Gamma)$ , set  $f_u(i) = G(u, i)$ . By (1.5) we know that  $f_u$  is a bounded linear functional and hence by the Riesz Representation theorem we know that there is a unique  $i_u \in F(E)$  such that  $f_u(i) = (i_u, i)$ . Define  $\tilde{I}(u) =$  $i_u$ . This is a well defined function by the uniqueness of the Riesz Representation theorem. Also, from the uniqueness of the Riesz representation theorem we see that  $\tilde{I}$  is linear. By the bound given by the Riesz representation theorem we know that  $\|\tilde{I}\| \leq 1$ . If  $u \in H(\Gamma)$ , then we just observe that  $G(u, i) = (Iu, i)_{F(E)}$ and hence by uniqueness of the Riesz Representation theorem we know that  $Iu = i_u = \tilde{I}(u)$ . The exact same analysis works for  $\tilde{\Phi}$ .

**Lemma 1.4.33.** We have that ker  $\widetilde{I} \subseteq W$ . Similarly we have that ker  $\widetilde{\Phi} \subseteq Y$ .

*Proof.* We can see this fairly easily just by looking at appropriate test functions. Suppose that  $u \in \ker \widetilde{I}$ , i.e. G(u, i) = 0 for all  $i \in F(E)$ . Let  $v_1$  and  $v_n$  be boundary vertices such that  $C = v_1 v_2 \cdots v_n$  is a path from  $v_1$  to  $v_n$ . We define the current function  $I_C$  to be just 1 along each edge  $v_j v_{j+1}$  and to be -1 along each edge  $v_i v_{i-1}$  and we observe that  $I_C$  is clearly in F(E) since the sum of the currents at any interior vertex is 0 and  $I_C$  is finitely supported. Hence  $G(u, I_C) = 0$ . But we observe that

$$0 = G(u, I_C) = 2(u(v_0) - u(v_n))$$

and hence u is constant on all boundary vertices and hence in W.

Now for the other claim, suppose that  $i \in \ker \tilde{\Phi}$ , i.e. that G(u, i) = 0 for all  $u \in Z(\Gamma)$ . The claim is also not that difficult. Just consider the set indicator function  $u = \chi_{\{v\}}$  where v is any vertex. We simply note that

$$0 = G(u, i) = 2 \sum_{v \sim v'} u(v) - u(v')$$

and hence *i* has zero current sums. This holds on the boundary, and hence  $i \in Y$ .

We have a stronger result for ker  $\tilde{I}$  and ker  $\tilde{\Phi}$ .

**Lemma 1.4.34.** We have that ker  $\widetilde{I} \subseteq H(\Gamma)^{\perp}$  and similarly ker  $\widetilde{\Phi} \subseteq L(E)^{\perp}$ .

*Proof.* We will show only the first claim since the argument for the second is identical. Let  $\phi \in \ker \widetilde{I}$ . Then  $G(\phi, i) = 0$  for all  $i \in F(E)$ . A fortiori we know that  $G(\phi, i) = 0$  for  $i \in L(E)$ . Using the fact that I is a bijection from L(E) to  $H(\Gamma)$ , we have that if  $u \in H(\Gamma)$  then  $0 = G(\phi, Iu) = (\phi, u)_{Z(\Gamma)}$ . Since this holds for all  $u \in H(\Gamma)$ , we have that  $\phi \in H(\Gamma)^{\perp}$ .

**Lemma 1.4.35.** The map  $\widetilde{I}$  maps  $H(\Gamma)^{\perp}$  into  $L(E)^{\perp}$  and similarly  $\widetilde{\Phi}$  maps  $L(E)^{\perp}$  in  $H(\Gamma)^{\perp}$ .

Proof. If  $\phi \in H(\Gamma)^{\perp}$  then  $(\phi, u)_{Z(\Gamma)} = 0$  for all  $u \in H(\Gamma)^{\perp}$ . Hence  $G(\phi, Iu) = 0$ . Since I maps  $H(\Gamma)$  onto L(E), we know that  $(\tilde{\Phi}, \phi i) = G(\phi, i) = 0$  for all  $i \in L(E)$  so  $\tilde{\Phi}\phi \in L(E)^{\perp}$ . The other claim follows from an identical argument.  $\Box$ 

**Remark 1.4.36.** The spaces  $H(\Gamma)$  and L(E) are *naturally* functionally dual to each other. We get this by noting that

$$(u_1, u_2)_{Z(\Gamma)} = G(I(u_1), u_2) = G(u_1, I(u_2)) = (I(u_1), I(u_2))_{F(E)}$$

and similarly

$$(i_1, i_2)_{F(E)} = G(\Phi(i_1), i_2) = G(i_1, \Phi(i_2)) = (\Phi(i_1), \Phi(i_2))_{Z(\Gamma)}$$

# 1.4.8 Dirichlet to Neumann and Neumann to Dirichlet maps

In a certain sense we have both Dirichlet-to-Neumann and Neumann-to-Dirichlet maps. We have maps which send vaild boundary data to the data from unique minimal functions. We have shown that these unique minimal functions are not in general the unique finite power  $\gamma$ -harmonic (resp. loop sum zero) voltage or current functions. Thus we don't in general have a Dirichlet-to-Neumann map or a Neumann to Dirichlet map in the traditional sense. We will create several definitions to make discussing these issues easier:

**Definition 1.4.37.** We define the map which sends valid boundary voltages to the boundary currents of the unique power minimizing function with those boundary values to be  $\Lambda_M$  and similarly we call  $H_M$  the function which sends valid boundary currents to the boundary voltages of the unique minimal power function with those currents. We call these maps the **minimal boundary data maps**. We call  $\Lambda_M$  the minimal Dirichlet-to-Neumann map and similarly for  $H_M$ .

**Definition 1.4.38.** If there exist maps which take valid boundary data to unique valid dual boundary data, we call the appropriate maps the **harmonic boundary data maps**.

The harmonic boundary data maps represent the tradition Neumann-to-Dirichlet maps as we have in the finite case. We now prove a theorem that gives conditions for when we could have such maps.

**Theorem 1.4.39.** If the function I maps  $M(\Gamma)$  into M(E), then there exists a well defined harmonic Neumann-to-Dirichlet map on finite power functions. If the function  $\Phi$  maps M(E) into  $M(\Gamma)$  then there is a well defined harmonic Dirichlet-to-Neumann map on finite power functions.

Proof. First suppose that I maps  $M(\Gamma)$  into M(E) and let  $i_0 \in L(E)$  be a current function that has zero boundary values, i.e.  $i_0 \in Y$ . Then  $(i, i_0) = 0$ for all  $i \in M(E)$  since  $M(E)^{\perp} = Y$ . But then, since I maps  $M(\Gamma)$  into M(E)and I is a bijection from  $H(\Gamma)$  to L(E), we know that for all  $u \in M(\Gamma)$ , there is an  $i \in M(E)$  such that  $\Phi(i) = u$ , and hence since  $\Phi$  is an isometry, we know that  $(i, i_0) = (u, \Phi(i)) = 0$ . Since this holds for all  $u \in M(\Gamma)$  we know that  $\Phi(i) \in W = M(\Gamma)^{\perp}$ . But, this means that if a finite power function has zero current on the boundary, then it has constant voltage, so that the harmonic Neumann-to-Dirichlet map is well defined. The other direction, i.e. showing that if  $\Phi$  maps M(E) into  $M(\Gamma)$  then there is a well defined harmonic Dirichlet-to-Neumann map on power functions, is proved just by switching the appropriate symbols.

**Corollary 1.4.40.** In general, we don't have that  $M(\Gamma)$  is mapped into M(E) by I or that M(E) is mapped into  $M(\Gamma)$  by  $\Phi$  since in general we don't have well defined harmonic boundary data maps.

**Remark 1.4.41.** If G is finite, then the spaces  $H(\Gamma)$  and  $M(\Gamma)$  correspond because there is a unique solution for given boundary data. Similarly L(E) and M(E) correspond.

We now analyze  $\Lambda_M$  and  $H_M$  a bit further. Perhaps our discussion is a bit out of order, but we now define the spaces of relevant boundary information.

**Definition 1.4.42.** Let  $\Omega(V)$  be the vector space of valid boundary voltages and let  $\Omega(E)$  denote the vector space of valid boundary currents (as sets we can view both as subspaces of the rather large space  $\mathbb{R}^{V}$ ).

The maps  $\Lambda_M$  and  $H_M$  are defined between these spaces.

# **1.5** The Space of Finitely Supported Functions

Here we analyze the space of finitely supported functions in both  $H(\Gamma)$  and L(E).

**Definition 1.5.1.** We define the space of finitely supported  $\gamma$ -harmonic functions to be  $H_f(\Gamma)$  and the space of finitely supported minimal functions to be  $M_f(\Gamma)$ . Similarly we define the space of finitely supported current loop sum zero current functions to be  $L_f(E)$  and we define the space of finitely supported minimal current functions to be  $M_f(E)$ .

**Lemma 1.5.2.** We have that  $H_f(\Gamma) = M_f(\Gamma)$  and  $L_f(E) = M_f(E)$ .

Proof. Obviously  $M_f(\Gamma) \subseteq H_f(\Gamma)$  and  $M_f(E) \subseteq L_f(E)$  by previous theorems. We now show that  $H_f(\Gamma) \subseteq M_f(\Gamma)$  Suppose  $v_0$  is an arbitrary vertex in Gand suppose that  $\phi$  is constant outside of  $G_n(v_0)$ . In particular, we know  $\phi$  is constant across any edge that is not in  $G_n(v_0)$ . We need to show that  $\phi \in W^{\perp}$ , so let  $w \in W$ . Considering the network  $\Gamma_{n+2}$  with graph defined by  $G_{n+2}(v_0)$ where the boundary vertices are  $(G_{n+2} \setminus G_{n+1}) \cup (\partial G \cap G_{n+2})$ . We observe that  $\phi$  is  $\gamma$ -harmonic on  $G_{n+2}(v_0)$  and hence is in  $W(\Gamma_{n+2})^{\perp}$ . There is a canonical map  $r_{n+2} : Z(\Gamma) \to Z(\Gamma_{n+2})$  defined by restricting a function on V(G) to a function on  $V(G_{n+2})$  (if  $u \in Z(\Gamma)$  denote its restriction to  $Z(\Gamma_{n+2})$  by  $u_{n+2}$ ). Noting that  $\phi$  has zero voltage difference across any edge that is not in  $G_{n+1}$ , so in particular for all the edges in  $G_{n+2} \setminus G_{n+1}$ , we note that

$$(\phi, w)_{Z(\Gamma)} = \sum_{V \times V} \gamma_{vv'}(\phi(v) - \phi(v'))(w(v) - w(v'))$$
$$= \sum_{G_{n+1} \times G_{n+1}} \gamma_{vv'}(\phi(v) - \phi(v'))(w(v) - w(v'))$$
$$= (\phi_{n+1}, w_{n+1})_{Z(\Gamma_{n+1})}.$$

But we notice that  $w_{n+1}$  is some constant c on  $\partial G \cap G_{n+1}$ . So we can extend  $w_{n+1}$  to be the function w' on  $G_{n+2}$  by setting w' to be c on all  $G_{n+2} \setminus G_{n+1}$ . We note then that  $w' \in W(G_{n+2})$  and furthermore

$$(\phi, w)_{Z(\Gamma)} = (\phi_{n+2}, w')_{G_{n+2}} = 0$$

since  $\phi_{n+2}$  is  $\gamma$ -harmonic and  $\gamma$ -harmonic functions correspond to minimal functions on finite graphs and  $w' \in W(G_{n+1})$ . Hence  $\phi \in W^{\perp}$  and hence  $\phi$  is minimal.

The proof that  $L_f(E) \subseteq M_f(E)$  is identical.  $\Box$ 

**Lemma 1.5.3.** In general we have that  $M_f(\Gamma)$  is mapped in  $M_f(E)$  by I but that  $M_f(E)$  is only mapped into  $M(\Gamma)$  by  $\Phi$ .

Proof. Clearly  $M_f(\Gamma)$  is mapped into L(E) by I. Furthermore, if  $u \in M_f(\Gamma)$  is finitely supported, then Iu will be finitely supported so  $Iu \in L_f(E)$ . By the previous theorem we know that  $L_f(E) = M_f(E)$ .

Now if  $i \in M_f(E)$ , suppose that i is supported on some  $G_n(v)$ . Then in particular, we know that  $\Phi i$  will be constant on each connected component of  $G_{n+1}(v) \setminus G_n(v)$ . We note that  $\Phi i$  is  $\gamma$ -harmonic on int G. Pick a representative of w such that w is zero on  $\partial G$ . If  $w \in W$  then we observe that

$$(\Phi i, w) = \sum_{V \times V} \gamma_{vv'} \left[ (\Phi i)(v) - (\Phi i)(v') \right] (w(v) - w(v'))$$
  
= 
$$\sum_{G_{n+1} \times G_{n+1}} \gamma_{vv'} \left[ (\Phi i)(v) - (\Phi i)(v') \right] (w(v) - w(v'))$$
  
= 
$$\sum_{v \in G_{n+1}} \sum_{v' \sim v} \gamma_{vv'} \left[ (\Phi i)(v) - (\Phi i)(v') \right] (w(v) - w(v')).$$
(1.6)

This expression is a finite sum and is linear in w. Let  $\delta_v$  denote the function which is 1 at v and 0 at all other vertices. We note that equation (1.6) can thus be rewritten as

$$\sum_{\rho \in (\operatorname{int} G) \cap G_n} w(\rho) \sum_{v \in G_{n+1}} \sum_{v' \sim v} \gamma_{vv'} \left[ (\Phi i)(v) - (\Phi i)(v') \right] (\delta_{\rho}(v) - \delta_{\rho}(v')),$$

but we notice that for a fixed  $\rho$  we have that

$$\sum_{v \in G_{n+1}} \sum_{v' \sim v} \gamma_{vv'} \left[ (\Phi i)(v) - (\Phi i)(v') \right] \left( \delta_{\rho}(v) - \delta_{\rho}(v') \right)$$
$$2 \sum_{v \sim \rho} \gamma_{v\rho} \left[ (\Phi i)(v) - (\Phi i)(v') \right]$$

which is zero since  $\Phi i$  is  $\gamma$ -harmonic on int G and  $\rho \in \text{int } G$ . Hence by linearity and the finiteness of the sums involved, we know that

$$(\Phi i, w) = \sum_{\rho \in (\inf G) \cap G_n} w(\rho) \sum_{v \in G_{n+1}} \sum_{v' \sim v} \gamma_{vv'} \left[ (\Phi i)(v) - (\Phi i)(v') \right] (\delta_{\rho}(v) - \delta_{\rho}(v')) = 0$$

so  $\Phi i \in W^{\perp} = M(\Gamma)$ , so we are done.

# 1.6 Half Planar and Dual Graphs

The natural analogue of circular planar graphs for infinite graphs turns out to be half planar:

**Definition 1.6.1.** We say an electrical network  $\Gamma$  is **half-planar** if  $\Gamma$  is embeddable in the upper half plane  $\overline{\mathbb{H}} \subseteq \mathbb{C}$  such that the boundary vertices are all on  $\mathbb{R}$  and the interior vertices are in  $\mathbb{H}$  (the strict upper half plane).

#### 1.6.1 Dual Networks

Just as in circular planar graphs we can construct a dual graph in the case of half planar networks. *A-priori* the dual graph doesn't need to have finite valence, so we will only consider primal graphs such that the dual graph has finite valence at every vertex. We can form an electrical network  $\Gamma^{\dagger}$  using  $G^{\dagger}$ by setting the conductance across any edge a'b' in the dual graph to be  $1/\gamma_{ab}$ where ab is the edge in G which crosses a'b'. Now given a  $\phi \in H(\Gamma)$ , then we wish to form the dual function  $\phi^{\dagger}$  on the dual function by requiring that if a'b'is the dual edge crossing ab and a'b' is the counterclockwise rotation of ab then

$$\phi^{\dagger}(b') - \phi^{\dagger}(a') = \gamma_{ab}(\phi(b) - \phi(a)).$$

There is nothing super sophisticated about this, but what we will do instead is form a current function on  $\Gamma^{\dagger}$  by setting  $i_{\phi}(b'a') = \phi(b) - \phi(a)$ . We have that  $i_{\phi}$ is  $\gamma$ -harmonic on int  $G^{\dagger}$  since  $\phi$  is a well defined voltage function and summing  $i_{\phi}$  at a vertex corresponds to summing voltage differences around a loop, which obviously yields zero. Similarly  $i_{\phi}$  has loop sum zero since it clearly has loop sum zero around any loop which bounds a single cell in the dual graph since summing around any such loop corresponds to checking that  $\phi$  is  $\gamma$ -harmonic. We don't go into these details too explicitly since they are identical to what happens in the finite case. We note that if  $\phi \in H(\Gamma)$  then  $P(i_{\phi})_{\Gamma^{\dagger}} = P(\phi)_{\Gamma}$  and hence  $i_{\phi} \in L(E^{\dagger})$ . We define the map  $D_{\Gamma}$  by

$$D_{\Gamma}: H(\Gamma) \to L(E^{\dagger}), \quad \phi \mapsto i_{\phi}$$

and we note that D is a linear isometry. As is easily verified, the composition of all the functions below in the obvious order is the identity function from  $H(\Gamma)$  to  $H(\Gamma)$ 

$$H(\Gamma) \xrightarrow{D_{\Gamma}} L(E^{\dagger}) \xrightarrow{\Phi_{\Gamma^{\dagger}}} H(\Gamma^{\dagger}) \xrightarrow{D_{\Gamma^{\dagger}}} L(E) \xrightarrow{\Phi_{\Gamma}} H(\Gamma)$$

is the identity.

#### 1.6.2 Voltage-Covoltage

Let  $\Omega(V)$  denote the vector space of valid boundary voltages and let  $\Omega(E)$ denote the space of valid boundary currents  $\gamma$ -harmonic functions. We observe that there is not in general well defined map from these space into any of the spaces of functions defined on all vertices, or all edges, but we do have a well defined map from  $\Omega(E)$  into the set of valid boundary covoltages  $\Theta$ , and this map is a bijection, as is easily verified. We will denote the map from  $\Omega(E)$  by  $\partial$ . We leave the details to the reader since they are just as they are in the finite case (see [4]). Hence we have the following diagram

$$\Omega(V) \xrightarrow{\Lambda_M} \Omega(E) \xrightarrow[iso]{} \Theta,$$

and similarly we have

$$\Theta \xrightarrow[\text{iso}]{\partial^{-1}} \Omega(E) \xrightarrow{H_M} \Omega(V)$$

**Corollary 1.6.2.** The minimal boundary data maps  $\Lambda_M$  and  $H_M$  contain equivalent information to the minimal voltage-covoltage and covoltage-voltage maps respectively (for half-planar graphs).

# Chapter 2

# **Inverse Problems**

# 2.1 Introduction

Here we present a proof that we can recover a class of infinite graphs that we will can critical half planar. Much of the work we present here is a generalization of work by Will Johnson, and a lot of theory, especially of convex sets of the medial graph, originated from him. The general strategy of our recovery process mirrors his process for finite graphs, though many of the proofs for the infinite case differ significantly from the finite case.

# 2.2 Preliminaries

### 2.2.1 Half-Planar Graphs

The class of graphs that we will attempt to recover is a subset of what we will call the half-planar graphs:

**Definition 2.2.1.** Let  $\Gamma = (\partial G, \operatorname{int} G, K)$  be a (possibly infinite) electrical network. We will say that  $\Gamma$  is **half planar** if there is an embedding of G into the closed upper half plane  $\overline{\mathbb{H}} \subseteq \mathbb{C}$  such that the following conditions are satisfied:

- 1. all vertices in  $\partial G$  are in  $\mathbb{R}$ ,
- 2. all vertices in int G are in  $\mathbb{H} = \overline{\mathbb{H}} \setminus \mathbb{R}$ ,
- 3. V(G) is a discrete set of  $\mathbb{C}$ ,

We can define the dual and medial graphs exactly as we would in the finite case, but things don't have to be as well behaved. Unfortunately for the extension, we do need things to be well behaved, so the class of graphs that we wish to recover becomes somewhat restrictive. We use as the infinite analogue of critical circular planar graphs a class of graphs which we call critical half planar graphs. Many of the below conditions are probably redundant, but we list them as assumptions to simplify the below exposition as much as possible.

**Definition 2.2.2.** A half planar graph G will be called critical with respect to a particular embedding if, with respect to that embedding, the following conditions are satisfied:

- 1. every vertex in the dual graph his finite valence,
- 2. there are no loops or self intersection of geodesics,
- 3. two geodesics intersect at at most one point,
- 4. the geodesics intersect the real axis at exactly two points. This implies that the geodesics are compact.
- 5. if  $K \subseteq \mathbb{C}$  is compact, then  $\{g : g[0,1] \cap K \neq \emptyset\}$  is finite.
- 6. the geodesics can be parametrized by smooth curves with nonvanishing derivative.
- 7. each geodesic cell is compact and has as boundary only finitely many geodesics.

### 2.2.2 Dirichlet and Neumann Data

We now recall some foundational material about infinite electrical networks, as was proven in the previous chapter. We encourage the reader to first read that chapter, since much of the material and terminology will be assumed. Firstly, in that paper we showed that there were well defined maps (both Dirichlet-to-Neumann and Neumann-to-Dirichlet) which take valid boundary data to the dual boundary data of a unique minimal function assuming both boundary data. We call these maps the minimal boundary data maps. We denote these maps by  $\Lambda_M$  and  $H_M$ . We still need to be careful though, since we don't not have harmonic boundary data maps in the traditional sense. I present several examples in that chapter. It might be a natural conjecture that for critical circular planar that we get well defined harmonic boundary data maps, but this turns out to be false, as we exhibit with the following example.

Example 2.2.3. We examine the "railroad track" graph in Figure 2.1.

We explicitly verify that this graph is critical half planar with the embedding below in Figure 2.2. Notice that technically this isn't embedded in a "half plan" but that doesn't matter since we can just conformally map the region onto the half plane.

Now we put conductivities on  $\Gamma$  in the way below and we immediately see that there is a  $\gamma$ -harmonic vertex function with finite power that is zero on the boundary and yields nonzero current flowing out of the boundary, so there is no Dirichlet-to-Neumann map in the traditional sense.



Figure 2.1: The "infinite railroad" graph examined in Example 2.2.3.



Figure 2.2: The medial and dual graphs of the "infinite railroad" 2.2.3. Note that primal vertices are notated with a solid circle and dual vertices are notated with an empty circle.

# 2.3 Convex and Closed Sets in the Medial Graph.

Here we develop the theory of convex and closed sets of the medial graph for infinite graphs, which was first developed by Will Johnson in [4] for finite critical circular planar. As in the finite case, this section will form the technical heart of the paper. Some results carry immediately over from the finite case, but most require substantially different proofs. Sometimes if c is a cell in the medial graph, we will regard c as a subset of  $\mathbb{C}$ , in which case we will always use c to refer to the Euclidean interior of the region c. Hopefully no confusion will be had.

### 2.3.1 Basic Definitions

**Definition 2.3.1.** Two cells in a medial graph M are adjacent if they share an edge. A connected set of cells X is one that is connected through adjacency.

**Definition 2.3.2.** If  $X \subseteq M$  we say that X has a corner at a vertex v in the medial graph if X contains exactly one of the cells that touch v. We see that X has an anticorner at v if X contains exactly three of the cells next to v (see Figure 2.3)



Figure 2.3: Corners and anticorners.

**Definition 2.3.3.** We will say X has a degenerate corner at a vertex v in the medial graph if X contains two cells which are diagonally opposite each other across v, but neither of the two cells adjacent to these two cells. See Figure 2.4.



Figure 2.4: Degenerate Corners.

### 2.3.2 Results from the Jordan Curve Theorem

We assume as part of the embedding that each geodesic g can be parametrized by a function  $g:[0,1] \to \overline{\mathbb{H}}$  such that g(1) and g(0) but the image of g doesn't intersect  $\mathbb{R}$  at any other points. We form the function  $\tilde{g}:[0,1] \to \overline{\mathbb{H}}$  such that on [0,1/2] we define  $\tilde{g}(t)$  to be g(2t) and on [1/2,1] we define  $\tilde{g}$  to parametrize the straight line segment from g(1) to g(0) on  $\mathbb{R}$ . Hence  $\tilde{g}$  is a Jordan curve and hence by the Jordan curve theorem  $\mathbb{C} \setminus \tilde{g}[0,1]$  consists of exactly two connected sets, one which is bounded and one which is unbounded. The bounded component is clearly a subset of  $\mathbb{H}$ . Furthermore, it's really easy to see that if U is the unbounded component then  $U \cap \mathbb{H}$  is also unbounded and connected. If g is any such geodesic, we will define B(g) to be the cells of the medial graph whose interiors lie in the bounded component of  $\mathbb{C} \setminus \tilde{g}[0,1] = \mathbb{H} \setminus \tilde{g}[0,1]$  and we will define U(g) to be the cells in the medial graph whose cells lie in the unbounded component of  $\mathbb{H} \setminus \tilde{g}[0,1]$ . Clearly  $U(g) \cup B(g) = M$ .

We first need some facts given by the piecewise Jordan curve theorem.

Lemma 2.3.4. Let  $\gamma : [0,1] \to \mathbb{C}$  be a piecewise continuously differentiable, closed, simple curve such with only finitely many nonsmooth points and suppose that  $\gamma'^+$  and  $\gamma'^-$  exist at every point where  $\gamma$  is not continuously differentiable. Suppose further that  $\gamma'$  is nonzero at the points where  $\gamma$  is differentiable. By possibly reversing the orientation of  $\gamma$ , we can ensure that  $i\gamma'$  (rotation counterclockwise) points into the bounded region bounded by  $\gamma$  and  $-i\gamma'$  (rotation clockwise) points into the unbounded component. Reversing the orientation of  $\gamma$  interchanges which vector points into which region.

*Proof.* I might fill in this proof later, but this is essentially taken out of Gamelin (page 250).  $\Box$ 

We now state some alternate versions of the Jordan curve theorem.

**Lemma 2.3.5.** Let  $\gamma : (0,1)$  be a simple curve in  $\mathbb{H}$  which escapes every compact set  $K \subseteq \mathbb{C}$ . Then  $\mathbb{H} \setminus \gamma(0,1)$  consists of exactly two components, at least one of which is unbounded.

Proof. Note that  $\gamma \to \infty$  as  $t \to 0$  and  $t \to 1$ . Hence we can extend  $\gamma$  to a simple closed curve in  $\hat{\mathbb{C}}$  (which is  $\mathbb{C}$  adjoined with a single point at  $\infty$ ) by setting  $\hat{\gamma}(0) = \hat{\gamma}(1) = \infty$ . The claim follows from the regular Jordan curve theorem for  $\hat{\mathbb{C}}$  along with the observation that one of the components that  $\hat{\mathbb{C}} \setminus \hat{\gamma}[0, 1]$  will contain the closed lower half plane, and neither component contains  $\infty$  and hence if  $U_1$  and  $U_2$  are the components of  $\hat{\mathbb{C}} \setminus \hat{\gamma}[0, 1]$  then  $U_1 \cap \mathbb{H}$  and  $U_2 \cap \mathbb{H}$  will both be connected.

**Lemma 2.3.6.** Let x be a cell in the medial graph such that all boundaries of x are geodesics. Then x is in B(g) for some geodesic g which travels adjacent to x.

Proof. Suppose to the contrary that  $x \in U(g)$  for all g which border x. We will use Lemma 2.3.4 in a strong way. Furthermore, we can assume our geodesics to be smooth except at the points where they intersect the real axis. Let the edges of x be (in clockwise order)  $e_1, e_2, \ldots, e_n$ . Let  $\gamma_0$  denote the piecewise smooth curve parametrizing these edges clockwise (i.e. points in the bounded region have winding number 1). Let the edge  $e_i$  of x correspond to geodesic  $g_i$ . Let  $x_{i,\ell}$  and  $x_{i,r}$  denote the points where  $g_i$  intersects the real axis, ordered so that  $x_{i,\ell} < x_{i,r}$ . We now need to consider parametrizations of  $g_i$ . Let  $\hat{g}_i$  first parametrize the simple closed piecewise smooth curve defined by having  $\hat{g}_i$  first parametrize the straight line from  $x_{i,\ell}$  to  $x_{i,r}$  and then parametrize the image of  $g_i$  starting at  $x_{i,r}$  and travelling to  $x_{i,\ell}$ . Since there are  $\epsilon_1, \epsilon_2$  such that  $[x_{i,\ell} + \epsilon_1, x_{i,r} - \epsilon_1] \times (0, \epsilon_2)$  is nonempty and a subset of  $B(g_i)$ , and that furthermore, the points in  $[x_{i,\ell} + \epsilon_1, x_{i,r} - \epsilon_1] \times (0, \epsilon_2)$  are clearly going to be to the left of  $g_i$  in the sense given in Lemma 2.3.4, we know that x is to the left of  $g_i$  (in sense of the previous lemma) iff  $x \in B(g_i)$ . Since x is not in B(g), we know that the points in x are to the right of every geodesic  $\hat{g}_i$ . We oriented  $\gamma_0$  so that the interior of x corresponded to being to the right of  $\gamma_0$ . Hence we know that the orientation on  $\gamma_0$  agrees with the orientation of  $\hat{g}_i$  on the intersection of their images (on the appropriate edges of x). We note that exactly one of  $x_{i+1,\ell}, x_{i+1,r}$  is in  $[x_{i,\ell}, x_{i,r}]$  since G has a critical embedding and so there are no lenses. If we start on a portion of  $g_{i+1}$  which is not in  $B(g_1)$  (such a portion of  $g_2$  and  $\gamma_0$  agree, we must have that  $g_2$  intersects  $g_1$  and enters into the interior of  $B(g_1)$ . Hence  $x_{i+1,\ell} \in [x_{i,\ell}, x_{i,r}]$  and  $x_{i+1,\ell} > x_{i,\ell}$ . We repeat this to get that  $x_{1,\ell} < x_{2,\ell} < x_{3,\ell} < \cdots < x_{n,\ell}$ . But since our ordering was cyclic (it didn't which edge on x we started with), we get that  $x_{n,\ell} < x_{1,\ell}$  and hence  $x_{1,\ell} < x_{1,\ell}$ , which is nonsense. Hence x must be in the bounded region of at least one of the geodesics which are adjacent to x.

**Lemma 2.3.7.** Let x be a cell in the medial graph such that x does not border  $\mathbb{R}$ . Then x is in U(g) for some g which neighbors x.

*Proof.* The proof is basically the same as the proof of Lemma 2.3.6. Suppose that x is in the bounded component bounded by every geodesic that boarders x. Let  $\gamma_0$  parametrize the boundary of x in a counterclockwise fashion, i.e. x is to the left of  $\gamma$ . Now let  $\hat{g}_i$  be defined as in the proof of Lemma 2.3.6. We note that the orientation of  $\hat{g}_i$  and  $\gamma_0$  agrees on the intersection of their images for each i. Since there can be no lenses in the medial graph, and using the assumptions about orientation, we see that  $x_{i+1,r} < x_{i,r}$  for all i. But because of the cyclic ordering of the edges of x, we see that  $x_{1,r} < x_{n,r} < x_{n-1,r} < \cdots < x_{1,r}$ , which is nonsense.

**Lemma 2.3.8.** Let x be a cell which borders  $\mathbb{R}$ . Then  $x \in B(g)$  for at least one of the geodesics which borders x.

*Proof.* The proof is almost identical to the proof of Lemma 2.3.6 with a slight modification. To do this, we will introduce *psuedogeodesics*, which are just closed intervals of  $\mathbb{R}$ , with the convention that if I is a closed interval, then  $U(I) = \mathbb{H}$  and  $B(I) = \emptyset$ . We note that if x borders  $\mathbb{R}$ , then x has a boundary which consists of psuedogeodesics and geodesics. If  $g_i$  is a psuedogeodesic, then we will define  $\hat{g}_i$  to be the linear parametrization of  $g_i$  from right to left. It is easily verified with this definition that the proof from Lemma 2.3.6 goes through essentially without change. We note that if g is a psuedogeodesic then  $x \notin B(g) = \emptyset$  and hence  $x \in B(g)$  for some real geodesic which borders x. We leave this verification to the reader.

Combining Lemma 2.3.6 and 2.3.8 we get

**Corollary 2.3.9.** Let x be a cell in the medial graph, then  $x \in B(g)$  for some geodesic g which borders x. A-fortiori  $x \in B(g)$  for some geodesic.

#### 2.3.3 Theorems about Minimal Numbers of Corners

**Definition 2.3.10.** We say that X is simply connected if X is connected and every cell in  $M \setminus X$  can be connected to a boundary cell in  $M \setminus X$  by a path of pairwise adjacent cells in  $M \setminus X$ 

**Definition 2.3.11.** A geodesic path is a path in the edges of the medial graph. A simple geodesic path is a path of geodesics which is a path of geodesics that is either never reaches a vertex twice, or a closed loop (i.e. doesn't reach a vertex twice except that the first and last vertices are the same).

**Definition 2.3.12.** A **psuedosimple** geodesic path is a geodesic path that either never reaches a single edge twice, or is a closed loop such that no edges are reached twice, except that the first and last edges are the same.

**Definition 2.3.13.** We define the boundary  $\partial X$  of a set  $X \subseteq M$  as the set of medial edges which boarder both a cell in X and a cell in  $M \setminus X$ .

**Definition 2.3.14.** If  $\gamma$  is a geodesic path which traverses part of the boundary of X. Then we say  $\gamma$  is **left-inwardly-oriented** (with respect to X) if every cell that is to the left of the edges of  $\gamma$  (with respect to the direction of traversal of *gamma*) is in X and every cell to the right is in  $M \setminus X$ .

**Definition 2.3.15.** If  $\gamma$  traverses a portion of  $\partial X$ , we define  $\gamma$  to be **component-following** if at every degenerate corner of X that  $\gamma$  reaches,  $\gamma$  follows the edge as shown in Figure 2.5. More precisely, if at every degenerate corner  $\gamma$  will turn so as to continue along the boundary of the same cell along which it came to the degenerate corner.



Figure 2.5: A curve  $\gamma$  is component-following if whenever it reaches a degenerate corner as above, it follows the dotted line.

**Lemma 2.3.16.** If  $X \subseteq M$  then  $\partial X$  can be parametrized by an edge disjoint family of left-inwardly-oriented, complement-following, pseudosimple geodesic paths or loops.

*Proof.* We will do this by induction. We will create a family of curves such that "to the right" corresponds to being in X. Suppose such a (possibly empty) family  $\mathcal{F}$  of pseudosimple, left-inwardly-oriented, and component-following geodesic

curves has been defined. Let  $e_1$  be geodesic edge which is in  $\partial X$  but not in any path. Now just extend e along  $\partial X$  in both directions. We now consider different possibilities of corners that we can reach. If we reach a node where two adjacent cells are in X but the other two are not, then we only have one choice. If we reach a corner then we also have no choice. Thus having started at e, we cannot reach an edge that's in another curve of  $\mathcal{F}$  by only passing along geodesics and turning at corners and anticorners since there have been no choices for any other path to make. The only case that we need to discuss is if the path reaches a degenerate corner. In that case, we have up to two choices, as illustrated in Figure 2.6.



Figure 2.6: The dotted lines represent the possible choices we have to pick at a degenerate corner.

If any of the the dashed edges in Figure 2.6 is in a curve in  $\mathcal{F}$ , then we pick the direction that preserves the component-following property. If any of the edges entering the node are are already in a path in  $\mathcal{F}$ , then we know that exactly two of them must be (since if three were already traversed then the fourth would have to be as well, contradicting the fact that we got to this node without traversing edges that are in  $\mathcal{F}$ ), and furthermore, since  $\mathcal{F}$  is assumed to have the component-following property, by looking at 2.6, we immediately see that the edge allowing our curve to have the component-following property cannot be the one in another path.

The only possibility for running into an edge that has already been defined is if we meet an edge of the same path we are defining, which is allowable. Also, these curves are pseudosimple since we assume that we just stop if we are ever forced to take an edge that we've already encountered. They are left-inwardlyoriented and component-following by our choice of travel at anticorners. Thus such a family of curves exists as stated.

**Lemma 2.3.17.** Let  $X \subseteq M$  be an arbitrary subset which doesn't contain any boundary cells of the medial graph. Then X is simply connected iff  $\partial X$  can be parametrized as simple geodesic path. The set X is finite and simply connected iff  $\partial X$  can be parametrized by a single closed simple geodesic path.

*Proof.* Let  $\mathcal{F}$  be an edge disjoint family of left-inwardly-oriented, complement following psuedosimple geodesics which traverse the entire boundary of X. Let  $\gamma$  be an arbitrary element of  $\mathcal{F}$ . We wish to show that  $\gamma$  is simple. Suppose  $\gamma$  is not simple. Let v be a medial node of self intersection of  $\gamma$ . Note that v must occur at a degenerate corner of X. We claim that  $\mathbb{C} \setminus \gamma$  must consist of at least three components. Clearly  $\mathbb{C} \setminus \gamma$  must have at least two components, since it is either a closed curve in  $\mathbb C$  or can be extended to a closed curve on  $\hat{\mathbb C}$  by setting the left and right endpoints to be  $\infty \in \mathbb{C}$ , and hence by basic complex analysis, the winding number is constant in connected components, but the winding number must change as we cross  $\gamma$  since  $\gamma$  doesn't repeat any edges. Suppose that  $\mathbb{C} \setminus \gamma$  consists of exactly two components. Since  $\gamma$  is componentfollowing, we now that the set of cells in X which are adjacent to an edge in  $\gamma$ must be connected. Hence the cells in X which are diagonally opposite must be in the same component. Since each of the cells in  $M \setminus X$  which touch v cannot be in the same component as the two cells in X since we cross  $\gamma$  to get to them (and hence the winding number with respect to  $\gamma$  would change), we know those two cells must be in the same connected component. Let  $x_1$  and  $x_2$ be arbitrary points in the two cells in  $M \setminus X$  which are adjacent to v (with  $x_1$ and  $x_2$  be difference cells). Since  $x_1$  and  $x_2$  are in the same component, there is a path (a continuous 1-1 function from [0,1] to  $\mathbb{C}$ ) from  $x_1$  to  $x_2$  which does not cross  $\gamma$  (and we can even assume that this path is linear in the two cells which contain  $x_1$  and  $x_2$ ). Concatenating this with a "nice" path that goes through v and starts at  $x_1$  and ends at  $x_2$  (what "nice" means is left to the reader) as in Figure 2.7, we get a simple closed curve, and hence we can apply the Jordan curve theorem to get that there are exactly two regions.



Figure 2.7: The path discussed above.

Clearly *exactly* one of the two cells in X which is adjacent to v must have an interior which is entirely in the bounded region and the other one must have an interior which is entirely contained in the unbounded region. But this contradicts the fact that the two cells in X which touch v are in the same component of  $\mathbb{C} \setminus \gamma$ . Now we now that all curves in  $\mathcal{F}$  must be simple. Now we claim that  $\mathcal{F}$  must consist of at most one path. Let  $\gamma$  be any of the paths in  $\mathcal{F}$ . Note that by the Jordan curve theorem  $\gamma$  will divide  $\mathbb{H}$  into exactly two connected regions. Note that exactly one of these regions can contain any cells in the boundary of the medial graph since X contains no cells of the medial graph and hence  $\partial M$  (the boundary cells of the medial graph) must be contained in the single component. Let Y consist of all of the cells whose interiors are in the other component. Obviously  $\partial Y = \operatorname{Im} \gamma$ . Suppose that X had some boundary edge e which was not contained in the image of  $\gamma$ . Then we note that either both cells must be in Y or both cells must be in  $M \setminus Y$ . If there is a cell from X in  $M \setminus Y$  then X cannot be connected to the boundary. We leave it to the reader to verify these last two assertions, but they are very clear.

For the last statement, if X is finite and simply connected and doesn't touch the boundary, then  $\partial X$  must be finite, and the only possibility for a simple geodesic curve to be finite is for it to be closed. If  $\gamma$  is a closed simple geodesic curve, then the fact that X is simply connected and finite follows from the Jordan curve theorem. The details are left to the reader.

**Corollary 2.3.18.** If X is simply connected, then X contains no degenerate corners.

**Lemma 2.3.19** (Lemma 5.4 of [4]). Let  $X \subseteq M$  be connected and finite and suppose X contains no boundary cells of the medial graph. Then X has at least three corners.

We leave it to the reader to verify that the proof from [4] carries over without change.

**Lemma 2.3.20.** Let X by a simply connected set and let g be a geodesic. Then  $X \cap B(g)$  and  $X \cap U(g)$  contain no degenerate corners.

*Proof.* If  $X \cap B(g)$  contained a degenerate corner, then X must too since none of the other two adjacent cells can be added when we pass to X since none of the geodesic passing through the vertex at this degenerate corner cannot be g since otherwise those two cells would not be diagonally adjacent in  $X \cap B(g)$ . The same argument holds for  $X \cap U(g)$ .

**Lemma 2.3.21.** Let X be simply connected and g a geodesic such that  $B(g) \cap X$  is nonempty. Then every component of  $B(g) \cap X$  has a corner which is also a corner of X.

*Proof.* Let C be a component of  $B(g) \cap X$ . If the segment of g that corresponds to the boundary of C is not connected, then we can fill in the components of  $B(g) \setminus X$  that touch the intermediate segments of g as shown in Figure 2.8.

This clearly does not reduce the number of corners. Call the resulting region  $\hat{C}$ . Clearly  $\hat{C}$  is finite and doesn't contain any of the boundary of M. Since  $\hat{C}$ 



Figure 2.8: Filling components along g.

is connected we know by Lemma 2.3.19 that  $\hat{C}$  has at least 3 corners. At most two of these cells can be adjacent to g since the segment of g that borders  $\hat{C}$  is connected. Hence one of these corners must occur inside of B(g) but not along g. Let c be such a corner. We note that clearly this corner is also a corner of C. Furthermore, by Lemma 2.3.20, we know that c is a corner of  $B(g) \cap X$ . Since is not adjacent to g, we know that c is not adjacent to any cells in U(g) and hence c is a corner of X. This holds for every component of  $B(g) \cap X$ .

We wish to show the infinite analogue of Lemma 2.3.19. It turns out this theorem is actually quite difficult to prove. It is relatively straightforward to show that an infinite simply connected subset which doesn't contain any boundary cells has at least one corner by just applying Lemma 2.3.19 to a subregion bounded by a geodesic. It is somewhat more complicated but still relatively straightforward to find a second corner by traveling in both directions from the corner we have already found. It turns out that finding the third corner is quite difficult, and the only proof that I could figure out actually implied that we have infinitely many corners.

**Lemma 2.3.22.** Let  $X \subseteq M$  be simply connected and infinite and suppose that X contains no boundary cells. Then X has infinitely many corners.

**Proof.** Let  $\gamma$  be a geodesic path which traverses the boundary of X. Suppose that X has no more than k corners (the case k = 0 is allowed). We observe that  $\gamma$  must be an infinite curve, since otherwise X would be bounded, contradicting our assumptions. After a certain point in traversing the boundary,  $\gamma$  can only travel along geodesics or meet anticorners and there must be infinitely many anticorners, since the geodesics are compact. Let  $c_1, \ldots, c_k$  be the corners of X. Since X is connected there are paths  $p_i$  of adjacent cells in X from  $c_i$  to  $c_1$ . By Lemma 2.3.9, we know that each cell c in the any of the paths  $p_i$  (which includes all  $c_i$ ) is in a region  $B(g_c)$  for some geodesics  $g_c$ . Define

$$R = \bigcup_{\{c: \exists p_j, c \in p_j\}} B(g_c)$$

and note that R is finite and contains all  $c_i$  for i = 1, ..., k and also R contains a path from each  $c_i$  to each  $c_j$ . Since R is finite, there can only be finitely many geodesics which pass through any portion of the interior of R or pass along the boundary of R. Since  $\gamma$  must be an infinite path which doesn't intersect itself (since otherwise X would be finite), we know that eventually one of the geodesics which has an edge on  $\gamma$  must not pass through R. Let g be such a geodesic. We note any component of g which  $\gamma$  traverses must start and end with an anticorner, since otherwise there would have to be a corner along g, which isn't possible since all the corners are contained in R and g doesn't border any cells in R. Hence  $X \cap B(g)$  and  $X \cap U(g)$  are both nonempty. Either B(g) or U(g)must contain all of R since g doesn't intersect the boundary of R. If  $R \subseteq U(g)$ then we apply Lemma 2.3.21 to show that  $B(g) \cap X$  contains a corner, which can't be any of the corners in R since  $R \subseteq U(g)$  and hence X has at least k + 1corners. Thus suppose that  $R \subseteq B(g)$ . Now we have to break into cases about which side of g corresponds to the bounded side with respect to the anticorners. Note that the anticorners on the segment of g that we are discussing must occur on the same side of g since otherwise one would be a corner.

Let the "sharp" side denote the side of g where the single cells of the anticorner are and let the "smooth" side of g refer to the side where the two adjacent cells of the anticorners are. See Figure 2.9 for a picture of what this means.



Figure 2.9: The smooth and sharp sides

If the "sharp" side of g is B(g), then we first claim that  $B(g) \cap X$  must consist of at least two components. To see this, let x and y be the two "sharp" cells in the anticorners adjacent to g as in Figure 2.10.

Suppose to the contrary that  $B(g) \cap X$  is connected and hence there is a path from x to y in  $B(g) \cap X$ . Let z be any cell along g between x and y which is not in X (such a cell obviously exists since otherwise there wouldn't be anticorners. But then there is clearly a loop of cells in X which bound z, and hence X is not simply connected. This is illustrated in Figure 2.11. Since all of the corners are connected by paths in R and  $R \subseteq B(g)$ , we know that  $R \cap B(g)$  all of the corners  $c_1, \ldots, c_k$  will be contained in a single component of  $X \cap B(g)$ . But since there is at least one other component of  $B(g) \cap X$ , we just apply Lemma 2.3.21 to get an additional corner in X, which can't be one of  $c_1, \ldots, c_k$  since it's in a different component of  $X \cap B(g)$  and hence X has k+1corners. Now suppose that the "smooth" side of g is bounded. If  $B(g) \cap X$  is



Figure 2.10: The cells x and y.

not connected, then one of the components must contain all of the corners in R as before and then another component will contain an extra corner of X. The last possible case is that  $B(g) \cap X$  is connected. Let  $e_1$  be any geodesic edge on the boundary of  $B(g) \cap X$  which is not along g and let  $e_2$  be any geodesic edge along g between the two anticorners x and y. We have the situation in Figure 2.12.

Since X is simply connected, by Lemma 2.3.17 there is a geodesic path  $\alpha$  from  $e_1$  to  $e_2$ . Since  $\alpha$  starts at  $e_1$  and the entire connected segment from x and y is part of the boundary of X, we know that  $\alpha$  must first travel along g and then turn at one of the anticorners x and y. If we truncate  $\alpha$  at the last edge before it either travels along g again or enters B(g) and then concatenate this curve with g to get back to the anticorner that  $\alpha$  turns at (see picture), then we form simple geodesic loop, which must bound a set of cells S in M which are not in B(g). Furthermore  $S \cap X$  is nonempty. We summarize with Figures 2.13 and 2.14.

By filling in holes along g if  $X \cap S \neq S$ , we can apply the same argument in Lemma 2.3.21 to see that S must have a corner that is also a corner of X. Since this corner is not in B(g), we know that it must be distinct from  $c_1, \ldots, c_k$ . Hence X has k + 1 corners.

**Corollary 2.3.23.** If  $X \subseteq M$  is connected and infinite and doesn't contain any boundary cells then X has infinitely many corners.

*Proof.* Just fill in all of the components of  $M \setminus X$  except the one that contains the boundary  $\partial M$ , which leaves a simply connected region.



Figure 2.11: A path in X which surrounds a cell in  $z \notin X$ .



Figure 2.12: The edges  $e_1$  and  $e_2$ .

# 2.3.4 Convex and Closed Sets

**Definition 2.3.24.** We define a half plane to be a set of the form B(g) or U(g) for some geodesic g.

**Definition 2.3.25.** A set  $X \subseteq M$  is closed if it has no anticorners.

Lemma 2.3.26. The intersection of arbitrarily many closed sets is closed.

*Proof.* Consider an arbitrary collection  $\{X_{\alpha}\}_{\alpha \in A}$  of closed sets. Let  $x_1, x_2, x_3, x_4$  be cells around a vertex v in the medial graph. If  $v_i \notin \bigcap_A X_{\alpha}$  then  $v_i \notin X_{\alpha}$  for some  $\alpha$ . Since  $X_{\alpha}$  is closed, another  $x_j \notin X_{\alpha}$  and hence  $x_j \notin \bigcap X_{\alpha}$  so  $\bigcap X_{\alpha}$  does not contain an anitcorner and hence is closed.

Lemma 2.3.27. Half planes are closed.



Figure 2.13: The curve  $\alpha$ . The curve  $\alpha$  is represented with a dashed line.



Figure 2.14: The region S. It is the region bounded by the dashed line.

Proof. Obvious.

Corollary 2.3.28. Convex sets are closed.

*Proof.* This follows from Lemma 2.3.27 and Lemma 2.3.26.

**Definition 2.3.29.** We define a set  $X \subseteq M$  to be convex if it is an intersection halfplanes.

**Theorem 2.3.30.** Every convex set X of cells is connected.

*Proof.* The proof caries over without alteration.

**Definition 2.3.31.** We define the closure of X to be the intersection of all closed sets which contain X.

**Lemma 2.3.32.** The closure of X can also be attained as a countable increasing chain of sets  $X = X_0 \subseteq X_1 \subseteq X_2 \subseteq \ldots$  such that each  $X_{i+1} \setminus X_i$  consists of at most one cell in the medial graph, and if  $x \in X_{i+1} \setminus X_i$  then x fits into an anticorner of  $X_i$ .

Proof. Let  $\mathcal{F}$  the collection of all supersets Y of X such that there is an at most countable chain of sets  $X = X_0 \subseteq X_1 \subseteq X_2 \subseteq \cdots$  such that  $Y = \bigcup_i X_i$  and each  $X_i$  differs from  $X_{i-1}$  by filling in a single corner. Let  $\mathcal{F}$  be ordered by set inclusion. We claim that  $\mathcal{F}$  contains upper bounds. Let  $Y_1 \subseteq Y_2 \subseteq \cdots$  all be sets in  $\mathcal{F}$ . Then we claim that  $\bigcup_i Y_i$  is in  $\mathcal{F}$ . This follows essentially from the proof that  $\mathbb{N} \times \mathbb{N}$  is countable. We'll use a picture to demonstrate what to do, but the argument is exactly the same as what one would canonically do to show that  $\mathbb{N} \times \mathbb{N}$  was countable. Let  $X = X_0^i \subseteq X_1^i \subseteq X_2^i \subseteq \cdots$  be an increasing chain of sets which increases at each step at most by filling in a single anticorner and whose union is  $Y_i$ . Take an increasing union of the sets  $X_j^i$  in the order presented in Figure 2.15.



Figure 2.15: Take an increasing union in this order.

If  $Z_j$  is the  $j^{\text{th}}$  set that we hit, let  $\widetilde{Z_j} = \bigcup_{i \leq j} Z_j$  and clearly  $Z_0 = X$  and  $Z_j \setminus Z_{j-1}$  consists of at most a single cell filled into an anticorner of  $Z_{j-1}$  and also  $Z \stackrel{\text{def}}{=} \bigcup_{\mathbb{N}} \widetilde{Z_j} = \bigcup_{\mathbb{N}} Y_i$ . So Z is an upper bound for all  $Y_i$  and  $Z \in \mathcal{F}$ . Finally we apply Zorn's lemma to get a maximal element  $M \in \mathcal{F}$ , which we claim is closed. Since M is written as a countable chain on increasing sets starting with X which differ by at most 1 cell added to a corner, we know that M must be in the closure of X. Hence  $M = \overline{X}$ .

Corollary 2.3.33. The closure of a connected set is connected.

*Proof.* This is a consequence of the previous lemma.

**Theorem 2.3.34.** If X is convex, then X is closed.

*Proof.* By Lemma 2.3.26, it is sufficient to show that half planes are closed, but this is obvious, since clearly you can't have an anticorner in a half plane.  $\Box$ 

**Definition 2.3.35.** If  $X \subseteq M$ , then let  $\widetilde{X}$  denote the intersection of all half planes containing X.

**Lemma 2.3.36.** If  $X \subseteq M$ , then  $\widetilde{X}$  is closed.

*Proof.* By Lemma 2.3.27 and Lemma 2.3.26 we know that X is an intersection of closed sets and is hence closed.

**Corollary 2.3.37.** If  $X \subseteq M$  then  $\overline{X} \subseteq \widetilde{X}$ .

Lemma 2.3.38. If X is connected and closed then X is simply connected.

*Proof.* We will use a similar argument to that found in [4]. Suppose X is not simply connected Consider all of the components of  $M \setminus X$ . Let S be a component which does not include any boundary cells. By Lemmas 2.3.23 and 2.3.19 we know that S must contain at least three corners  $(v_i, x_i)$  (using Will's notation). The two cells adjacent to  $x_i$  must be in X since  $(v_i, x_i)$  are corners, but the cell across from  $x_i$  must be in  $M \setminus X$  since if the cell were in X, then X would contain an anticorner, which would contradict the fact that X is closed. Thus we can form an adjacency multigraph A for connected components of  $M \setminus X$  by having each component be a vertex and connecting to vertices iff they have corners which are diagonally opposite to each other at the same vertex. By Lemma 2.3.19 and Lemma 2.3.23, we know that each component of  $M \setminus X$ which does not intersect the boundary has at least three edges connected to it. Now pick an interior component  $A_0$  of  $M \setminus X$  arbitrarily. We now we will define a path inductively in A which extends in both directions from  $A_0$ . Since each interior component has degree at least 2, we know that if our path ends in an interior vertex, we can always extend farther in that direction. If we get a cycle, then we know X is disconnected (by applying the Jordan Curve theorem). There are four possibilities:

- 1. the path forms a cycle,
- 2. the path extends indefinitely in both directions,
- 3. the path extends indefinitely in one direction, but reaches the boundary in the other,
- 4. the path reaches the boundary in both directions.

In all cases, we get that X is disconnected and hence there can be no connected components of  $M \setminus X$  which don't intersect the boundary.

We now diverge from the treatment of closed and compact sets that is found in [4] and introduce some new lemmas and results.

**Definition 2.3.39.** Let  $x = x_0, x_1, \ldots, x_n = y$  be a path of adjacent cells in the medial graph such that no cell is repeated. Let  $p_0, \ldots, p_n$  be arbitrary points such that  $p_i \in x_i^{\circ}$  (the interior of  $x_i$ ) for all *i*. Let  $\gamma : [0, 1] \to \mathbb{C}$  be a curve such that the following conditions are satisfied:

1. 
$$\gamma(k/n) = p_k$$
 for  $k = 0, ..., n$ ,

- 2.  $\gamma[k/n, (k+1)/n] \subseteq \operatorname{int} \overline{x_k \cup x_{k+1}}$
- 3.  $\gamma$  is piecewise smooth.
- 4.  $\gamma$  does not intersect itself.

Then we will call  $\gamma$  a continuous curve associated with the path  $x_0, \ldots, x_n$ .

**Lemma 2.3.40** (Filling Lemma). Let  $X \subseteq M$  be a closed connected set. Let  $x_1, x_2$  be two cells on the same side of g and let  $y_1 = x_1, y_2, \ldots, y_n = x_2$  be a path of adjacent cells from  $x_1$  to  $x_2$  which lies entirely on the same side of g as  $x_1$  and  $x_2$ . Then there is a region R of medial cells which is bounded by g and the path  $y_1, \ldots, y_n$  and R consists of a finite union of simple connected regions and furthermore  $R \subseteq X$ . In particular all of the cells along g between  $x_1$  and  $x_2$  are in X.

*Proof.* Let  $\gamma : [0,1] \to \bigcup_i \overline{y_i}$  be a continuous curve associated with the medial path  $y_1, \ldots, y_n$ . Extend  $\gamma$  to get a function  $\hat{\gamma} : [-1,2]$  such that

- 1.  $\hat{\gamma}(t) = \gamma(t)$  for  $t \in [0, 1];$
- 2.  $\hat{\gamma}(t) \in \overline{y_1}$  for  $t \in [-1, 0]$ ;
- 3.  $\hat{\gamma}(t) \in \overline{y_n}$  for  $t \in [1, 2]$ ;
- 4.  $\hat{\gamma}(-1) \in g \cap \overline{y_1}$  and  $\hat{\gamma}(2) \in g \cap \overline{y_n}$
- 5.  $\hat{\gamma}$  has no self intersections.
- 6.  $\hat{\gamma}$  intersects the geodesic g at exactly two locations.

Under these assumptions we can extend and reparametrize  $\hat{\gamma}$  to a function  $\tilde{\gamma}$  that is a Jordan curve by letting  $\tilde{\gamma}$  traverse the geodesic arc between  $\hat{\gamma}(-1)$  and  $\hat{\gamma}(2)$ . We will assume that  $\tilde{\gamma}$  is parametrized on the interval [0, 1]. We summarize in Figure 2.16.

Thus R exists as stated by noting that R is just the set of medial cells in  $M \setminus X$  which are in the region bounded by  $\tilde{\gamma}$ . It may not be true that R is connected, but it is sufficient to show that the connected component closest to  $x_1$  (along g) is in X, so without loss of generality we may assume that  $x_2$  is the first element of the path  $y_1, y_2, \ldots, y_n$  other than  $x_1$  which touches g. Let  $c_1, c_2, \ldots, c_m$  denote the cells between  $x_1$  and  $x_2$  along g such that  $x_1 = c_1$  and  $x_2 = c_m$ . Without loss of generality we may assume that  $\{c_2, \ldots, c_{m-1}\} \cap \{y_1, \ldots, y_n\}$  is empty, i.e. that  $c_m$  is the closest cell to  $x_1$  in  $\{y_2, \ldots, y_n\}$  which is along g and in the same direction as  $x_2$ . Our goal is to show that R is empty. We draw a picture to summarize the situation in Figure 2.17.

We first show that R is connected. Let  $\mathcal{R}$  denote the set of connected components of R. Define an adjacency graph A on  $\mathcal{R}$  by setting two components adjacent if they share a degenerate corner. By Lemma 2.3.19 each component



Figure 2.16: The curve  $\tilde{\gamma}$ .



Figure 2.17: The cells  $c_1, \dots, c_m$  and R. Then entire white region bounded by the grey loop and g is R.

of R has at least 3 corners. Since by assumption there is at most one component which has cells on g, and the geodesic segment which borders R must be connected (since all of the cells in the path  $y_1, \ldots, y_n$  are on one side of g), we know at most two cells from the component along g can correspond to corners along g, and that all other components have no corners along g. Hence at least one must correspond to an actual to either an anticorner of X or a degenerate corner of R. But X has no anticorners, and hence each connected component of S must have a degenerate corner with another connected component of S, and hence A is connected. But furthermore, we know that at most one connected component can have cells which are adjacent to g, and hence every vertex of A except for possibly a single vertex has valence 3 or greater. By forming a path in A by starting at the component along g and travelling in any manner such that we don't backtrack, we must eventually form a loop, thus forming a disconnection of X, a contradiction. Hence R must be empty.

**Lemma 2.3.41.** Suppose X is a closed connected subset of M, then  $\widetilde{X} = X$ .

*Proof.* We will show that there are no cells in  $\widetilde{X} \setminus X$  which are adjacent to X. Since  $\widetilde{X}$  is connected and contains X, this would imply that  $\widetilde{X} = X$ . We use the previous lemma. Suppose that  $c \in \widetilde{X} \setminus X$  is adjacent to X. We will show that  $c \in X$ . Let c be adjacent to  $x_0 \in X$  and let g denote the geodesic that travels between c and  $x_0$ . Since  $c \in \widetilde{X}$ , there must be some medial cell y in Xthat is on the same side of g as c. Since X is connected there is a path  $z_1, \ldots, z_n$ of adjacent cells in X such that  $z_1 = x$  and  $z_n = y$ . Let k denote the last index such that all  $z_i$  are on the same side of g as x for all  $i = 1, \ldots, k$ . Notice that  $z_{k+1}$  is on the same side of g as c. We note that the path  $z_1, \ldots, z_k$  is a loop only on one side of a geodesic with  $z_1$  and  $z_k$  along g. We summarize with the situation in Figure 2.18



Figure 2.18: The situation: cells  $y_1, \ldots, y_k, y_{k+1}$  and c.

Thus we are in a position to apply Lemma 2.3.40, and hence we know that all the cells along g between  $z_1$  and  $z_k$  on the same side of g as the path  $z_i$  must be in  $\overline{X}$ . But then since  $z_{k+1}$  is on the opposite side of g, and is adjacent to  $z_k$ , which is adjacent to a cell x' in X which is diagonally across from  $z_{k+1}$  (which is guaranteed to be in X by the filling lemma), we know that the fourth cell adjacent to  $z_{k+1}$  and x' is in X. By induction, all of the cells along g on the same side as  $z_{k+1}$  which are between  $z_{k+1}$  and c must be in X. Hence c must be in X, so we are done.

**Corollary 2.3.42.** If X is connected, then  $\overline{X} = \widetilde{X}$ .

*Proof.* We have that  $\widetilde{X} \supseteq \overline{X}$  since  $\widetilde{X}$  is closed and contains X. On the other hand, we have by the previous lemma that  $\widetilde{X} \subseteq \overline{\overline{X}} = \overline{X}$  and hence combining these two results yields the desired equality.

# 2.4 An Important Computation

**Lemma 2.4.1.** If X and Y are subsets of M then we have  $\overline{X} \cup \overline{Y} \subseteq \overline{X \cup Y}$  and  $\widetilde{X} \cup \widetilde{Y} \subseteq \overline{X \cup Y}$ .

*Proof.* Left to reader.

**Lemma 2.4.2.** Let g be a geodesic and let  $X_1$  consist of all of the boundary cells in the medial graph to the left of B(g) and let  $X_2$  consist of all of the boundary cells to the right of B(g). Then  $\overline{X_1 \cup X_2} = U(g)$ .

*Proof.* Clearly  $X_1 \cup X_2 \subseteq U(g)$  since U(g) is a half plane which contains both  $X_1$  and  $X_2$ . Note by Lemma 2.3.42 that since  $X_1$  and  $X_2$  are closed we have  $\overline{X_1} = X_1$  and  $\overline{X_2} = X_2$ . Let  $R = \overline{X_1 \cup X_2} \cup B(g)$  and consider  $S = M \setminus R$ . By Lemma 2.3.19 and Corollary 2.3.23 we know that each component of S has at least three corners. Now we observe that by the Filling lemma we know that the component of g which is not adjacent to any cells in  $\overline{X_1 \cup X_2}$  is connected. Now if  $S_0$  is the component of S which is adjacent to this connected arc of g, then by the previous lemmas we know that  $S_0$  has at least three corners, and at most two of them can be adjacent to g. The third corner cannot be an anticorner of  $X_1 \cup X_2$  since  $X_1 \cup X_2$  is closed. Hence the diagonal cell must be another component of S, which must have at least three corners. Furthermore, every component of S other than  $S_0$  must have at least three corners, all of which must be degenerate corners. Form an adjacency graph A on the components of S other than  $S_0$  and note that any loop or infinite path in A will create portion of  $\overline{X_1 \cup X_2}$  which cannot be connected to the boundary with a path of adjacent cells in  $\overline{X_1 \cup X_2}$ . Hence S must be empty, so  $\overline{X_1 \cup X_2} = M \setminus B(g)$ . 

# 2.5 Extending Consistent Functions

We will now reference some work done by Will Johnson in [4] that will carry over essentially without change to the infinite case. These initial theorems and definitions are essentially identical to what is presented in [4].

**Definition 2.5.1.** Let  $X \subseteq M$ . Let a be a cell in the medial graph. Then if  $X' = X \cup \{a\}$  we say that X' is a **simple** extension of X if a and three corners in X meet at an anticorner of X. We say X' is a **nice** simple extension if a touches exactly one anticorner of X (i.e. is adjacent to exactly two cells in X). If X'' is obtained from X by a series of simple extensions then X'' is an extension of X. If X'' is obtained by a sequence of simple nice extensions, then X'' is a nice extension of X.

**Theorem 2.5.2** ([4]). Let M be a finite critical medial graph and let  $X \subseteq M$  be convex such that  $X \cap \partial M \neq \emptyset$ . Then there is some set of cells  $S \subseteq \partial M \setminus X$  such that the entire medial graph is a nice extension of  $X \cup S$ .

Another result of Will Johnson's work that we get trivially is that

**Remark 2.5.3.** If c is a cell in  $\partial M \setminus X$  which is adjacent to X, then we can pick S in Theorem 2.5.2 such that  $c \in S$ .

We leave it to the reader to verify the last statement, but it follows trivially. We wish to apply Theorem 2.5.2 to sets of the form B(g) for geodesics in a medial graph. Though the theorem seems like it should obviously carry over in some form, we will provide justification, and then leave some of the details to the reader. The idea is that if X = U(g) for some geodesic g, then we want to be able to a single cell a from  $\partial M$  which borders X so that  $\overline{X \cup \{a\}}$  is a simple extension of X, but the theorem doesn't quite apply. To make it apply, we will construct a new medial graph, M', which has B(g) embedded so that

- 1. all geodesics in B(g) are geodesics in M';
- 2. the geodesic g is a geodesic in M'
- 3.  $M' \setminus B(g)$  is a closed and connected set of cells
- 4. M' is a critical medial graph.

If the above four conditions are satisfied, then we can just apply 2.5.2 to see that all M is a simple extension of X = U(g) since any simple extensions of a superset of X in M corresponds to a simple extension of a superset of  $M' \setminus B(g)$ in the obvious fashion.

We now need to construct such a medial graph M'. Thus turns out to actually be really easy. First, conformally map B(q) (as a closed subset of  $\mathbb{C}$ ) onto the closed unit disc. This is possible via the Riemann Mapping Theorem (which ensures that such a conformal map exists between the interiors of those regions) and Catheodory's Theorem (which states that since the boundaries of both regions are Jordan curves, the conformal map shown to exist by the Riemann mapping theorem extends continuously to the boundary). Now smooth each of the geodesics near the boundary so that if g is a geodesic then (if we view g as a function from [0,1] to B(g) (reparametrizing if necessary)) then  $\lim_{t\to 1^-} g'(t)$  and  $\lim_{t\to 0^+} g'(t)$  both exist. We can do this by simply replacing g with a linear function sufficiently close to the boundary such that we don't get any additional intersections of regions. Now let  $z_1$  and  $z_2$  be the two points on  $\partial \mathbb{D}$  which correspond to the points of intersection of g and  $\mathbb{R}$  in our original medial graph and let  $R_1$  and  $R_2$  denote the two arcs of  $\partial \mathbb{D}$  between  $z_1$  and  $z_2$ . Without loss of generality let  $R_1$  denote the arc corresponding to the geodesic g. Extend two rays segments from  $z_1$  and  $z_2$  radially. Call these line segments  $\ell_1$  and  $\ell_2$ . Now extend each geodesic which intersects  $R_1$  along a straight line corresponding to the limit of its derivative as it approaches  $R_1$ . Now using a basic compactness argument, we can pick a  $\rho > 1$  such that none of the extended geodesics or  $\ell_1$  or  $\ell_2$  intersect in the closed ball of radius  $\rho$ . Let  $x_1$  and  $x_2$  denote two points along  $\ell_1$  and  $\ell_2$  of norm  $\rho$  and let R denote the circular arc  $\rho R_1$ . Let H be the region bounded by  $R_2$ , the two line segments  $[z_1, x_1], [z_2, x_2]$  and the arc R. Clearly the extended geodesics will divide H into regions, which can be interpreted as medial graph cells for a medial graph M' in the obvious way. If so desired we can map the region H to the disc and then smooth the geodesics at the boundary as before. This procedure is summarized in Figure 2.19



Figure 2.19: An outline of the method of extending B(g) so that B(g) is the complement of a convex subset containing a connected subset of the boundary of a critical circular planar medial graph.

Obviously the medial graph M' satisfies the stated properties. Using Theorem 2.5.2 and the above observations, we get that:

**Theorem 2.5.4.** If M is an critical half planar graph and X = U(g), then there is a finite set of boundary cells  $S \subseteq \partial M$  such that M is a simple extension of  $X \cup S$ . Furthermore, if  $c \in B(g) \cap \partial M$  and c borders g, then we can pick Sso that  $c \in S$ .

# 2.6 Recovery

We are now nearly complete with most of our technical work and we need only to prove some basic facts which will allow for recovery. We first will show how to recover boundary spikes and boundary-to-boundary edges, which will turn out to be sufficient.

### 2.6.1 Recovering Boundary Spikes

A boundary spike is a pair of vertices  $(v_{\partial}, v_{\text{int}})$  such that  $v_{\partial} \in \partial G$  and  $v_{\text{int}} \in$ int G and there is an edge between  $v_{\partial}$  and  $v_{\text{int}}$  but there are no other edges connected to  $v_{\partial}$ . Throughout we assume that boundary spikes are never boundaryto-boundary edges, i.e. every boundary spike consists of a boundary vertex connected to exactly one interior vertex and no other boundary vertices.

Given an electrical network  $\Gamma$  with a boundary spike, we can form the network  $\Gamma_0$  where we remove the given boundary spike. There is an natural map from  $\Gamma_0$  into  $\Gamma$  which maps a vertex in  $\Gamma_0$  to the corresponding vertex in  $\Gamma$ . We will call this map  $\alpha : \Gamma_0 \to \Gamma$ . We have a simple but important lemma:

**Lemma 2.6.1.** Let  $\Gamma$  be an infinite electrical network with a boundary spike  $(v_{\partial}, v_{\text{int}})$  and let  $\Gamma_0$  be the electrical network corresponding to contracting  $(v_{\partial}, v_{\text{int}})$ . Let  $\phi \in Z(\Gamma)$ . If  $\phi \in M(\Gamma)$  then  $\phi_0 \stackrel{\text{def}}{=} \phi \circ \alpha \in M(\Gamma_0)$ . If  $\phi \in H(\Gamma)$  and  $\phi_0 \in M(\Gamma)$  then  $\phi \in M(\Gamma)$ .

**Proof.** Let W be the subset of  $Z(\Gamma)$  consisting of functions which are constant on the boundary of  $\Gamma$  and let  $W_0$  be the subset of  $Z(\Gamma_0)$  consisting of functions which are constant on the boundary of  $\Gamma_0$ . Notice that there is an obvious embedding of  $W_0$  into W by sending a function u which takes value c on the boundary of  $\Gamma_0$  to the function which is u on Im  $\alpha$  and takes value c on  $v_{\partial}$ . If  $u \in W_0$  let  $\tilde{u}$  denote the element of W as described.

Now to proceed with the proof of the lemma, suppose that  $\phi \in M(\Gamma) = W^{\perp}$ , then we wish to show that  $\phi_0 \in W_0^{\perp}$ . But to do this, we just note that if  $u \in W_0$ then  $\tilde{u} \in W$  and hence

$$\begin{aligned} (\phi_0, u)_{Z(\Gamma_0)} &= \sum_{V(G) \setminus \{v_\partial\} \times V(G) \setminus \{v_\partial\}} \gamma_{vv'}(\phi_0(v) - \phi_0(v'))(u(v) - u(v')) \\ &= 2\gamma_{v_\partial v_{\text{int}}} \left(\phi(v_\partial) - v_{\text{int}}\right) (\widetilde{u}(v_\partial) - \widetilde{u}(v_{\text{int}})) \\ &+ \sum_{V(G) \setminus \{v_\partial\} \times V(G) \setminus \{v_\partial\}} \gamma_{vv'}(\phi_(v) - \phi_(v'))(u(v) - u(v')) \\ &= (\phi, \widetilde{u})_{Z(\Gamma)} \end{aligned}$$

so that  $\phi_0 \in W_0^{\perp}$ . Now to prove the other direction, suppose that  $\phi_0 \in W_0^{\perp}$ . Let  $u \in W$ , without loss of generality, assume that u is 0 on  $\partial G$ . Let  $u_0$  denote  $u \circ \alpha$ . Note that  $u = \widetilde{u_0} + \chi_{v_{\text{int}}} u(v_{\text{int}})$  where  $\chi_{v_{\text{int}}}$  denotes the indicator function on the set  $\{v_{\text{int}}\}$ . Hence

$$(u,\phi)_{Z(\Gamma)} = (\widetilde{u_0} + \chi_{v_{\text{int}}} u(v_{\text{int}}), \phi)_{Z(\Gamma)}$$
  
=  $(\widetilde{u_0}, \phi)_{Z(\Gamma)} + (\chi_{v_{\text{int}}} u(v_{\text{int}}), \phi)_{Z(\Gamma)}$ 

By the previous computation, we know that  $(\widetilde{u_0}, \phi)_{Z(\Gamma)} = (u_0, \phi_0)_{Z(\Gamma_0)}$ which is zero by assumption. On the other hand, we know that  $(\chi_{v_{\text{int}}} u(v_{\text{int}}), \phi)_{Z(\Gamma)}$ is zero since it is exactly the current entering the vertex  $v_{\text{int}}$  and  $\phi \in H(\Gamma)$ . Hence  $\phi \in W^{\perp} = M(\Gamma)$  so we are done.

**Lemma 2.6.2.** Given the minimal Dirichlet-to-Neumann map  $\Lambda_M(\Gamma)$  and a boundary spike with given conductance  $\gamma_{v_\partial v_{int}}$ , we can find the minimal Dirichlet-to-Neumann map  $\Lambda_M(\Gamma_0)$  for the network  $\Gamma_0$  with the boundary spike contracted.

Proof. Let  $G_0$  denote the graph of  $\Gamma_0$  (the contracted network) and let  $\phi$ :  $\partial G_0 \to \mathbb{R}$  be a function such that there exists a  $\tilde{\phi}: V(G_0) \to \mathbb{R}$  of finite power such that  $\phi = \hat{\phi}|_{\partial G_0}$ . Now define  $\hat{\phi}$  to be  $\tilde{\phi}$  on  $V(G_0)$  and define  $\hat{\phi}$  to be the unique  $\gamma$ -harmonic extension of  $\tilde{\phi}$  to the boundary spike. Clearly this exists and is unique. Furthermore,  $\hat{\phi}$  is also clearly finite power. By Lemma 2.6.1, we know that  $\tilde{\phi}$  is of minimal power for its boundary voltages. We should note that  $\phi: \partial G_0 \to \mathbb{R}$  and  $\Lambda_M(G)$  uniquely determines  $\tilde{\phi}(v_\partial)$  since

$$\Lambda_M(G)(\hat{\phi}|_{\partial G})(v_{\partial}) = \gamma_{v_{\partial}v_{\rm int}} \left(\phi(v_{\rm int}) - \hat{\phi}(v_{\partial})\right),$$

and  $\gamma_{v_{\partial}v_{\text{int}}}$  is nonzero so we can just solve for  $\hat{\phi}(v_{\partial})$  in terms of known quantities. By  $\gamma$ -harmonicity we know that

$$\Lambda_M(G_0)(\phi)(v_{\text{int}}) = \Lambda_M(G)(\phi|_{\partial G}).$$

<u></u>

Furthermore,  $\Lambda_M(G)(\widetilde{\phi}|_{\partial G})(v) = \Lambda_M(G_0)(\phi)(v)$  for  $v \in \partial G$  such that  $v \neq v_{\text{int}}$ . Hence  $\Lambda_M(G)$  and the conductivity  $\gamma_{v_{\text{int}} v_{\partial}}$  uniquely determine  $\Lambda_M(G_0)$ .

**Lemma 2.6.3.** Given a critical half planar electrical network  $\Gamma$  with a boundary spike  $(v_{\partial}, v_{\text{int}})$ , the minimal Dirichlet-to-Neumann map  $\Lambda_M$  uniquely determines the conductivity  $\gamma_{v_{\partial}v_{\text{int}}}$ .

*Proof.* Let  $x_{\partial}$  denote the medial graph cell corresponding to  $v_{\partial}$ . We note that  $x_{\partial}$  is a geodesic triangle, i.e. there are two geodesics which bound  $x_{\partial}$  and the other edge of  $x_{\partial}$  is an interval of  $\mathbb{R}$ . By Lemma 2.3.9 we know that  $x_{\partial} \in B(g)$  for one of the two geodesics which is borders  $x_{\partial}$ . We will progressively define the function  $(\phi, \psi)$  where  $\phi$  is a voltage function on  $\Gamma$  and  $\psi$  is a covoltage function (on  $\Gamma^{\dagger}$ ) such that

$$\psi = (\Phi_{\Gamma^{\dagger}} \circ D_{\Gamma})(\phi).$$

Firstly define  $\phi$  and  $\psi$  to be 0 on  $\partial M \cap U(g)$ . By Lemma 2.4.2 we know that  $\phi$  and  $\psi$  are zero on U(g). By Lemma 2.5.4, we know that we can pick an  $S \subseteq \partial M$  such that  $x_{\partial} \in S$  and M is a simple extension of  $X \cup S$ . Define  $\phi$  to be 1 on  $x_{\partial}$  and specify  $\phi$  and  $\psi$  arbitrarily on the other cells of S. Just as in the finite case, we know that under these conditions we can extend  $\phi$  and  $\psi$  to be defined

on all of V and  $V^\dagger$  such that  $\phi$  is  $\gamma\text{-harmonic on int}\,G$  and  $\psi$  is  $\gamma^\dagger\text{-harmonic on int}\,G^\dagger$  and

$$\psi = (\Phi_{\Gamma^{\dagger}} \circ D_{\Gamma})(\phi)$$

(we leave these details to the reader, but it is identical to the finite case, so the interested reader can read [4]). Thus we can find a  $\gamma$ -harmonic function  $\phi: V \to \mathbb{R}$  such that  $\phi(v_{\partial}) = 1$  and  $\phi(v) = 0$  for  $v \in U(g) \cap \partial M$ . Further we have that if  $\psi = (\Phi_{\Gamma^{\dagger}} \circ D_{\Gamma})(\phi)$  then  $\psi$  is (up to a constant) 0 on  $U(g) \cap \partial M$ . By Lemma 2.4.2 we know that any functions  $\phi$  and  $\psi$  which satisfy those conditions will also satisfy  $\phi(v_{\text{int}}) = 0$ . We note that  $\phi$  is of finite power since it is finitely supported, and furthermore, since it is finitely supported, by Lemma 1.5.2 we know that  $\phi \in M(\Gamma)$ . Thus we can find a  $\phi$  and  $\psi$  satisfying the above conditions such that (using the notation from the previous chapter) that

$$(\partial \circ \Lambda_M)(\phi|_{\partial G}) = \psi|_{\partial G}$$

Now we are in the same situation as in the finite case, and we know that

$$\Lambda_M(\phi|_{\partial G})(v_{\partial}) = \gamma_{v_{\text{int}} v_{\partial}} \cdot (1-0)$$

which immediately gives us  $\gamma_{v_{\text{int}} v_{\partial}}$ .

## 2.6.2 Recovering Boundary-to-Boundary Edges

We recover boundary-to-boundary edges in a very similar fashion to how we recovered boundary spikes. We define a boundary-to-boundary edge to be an edge  $v_1v_2 \in E(G)$  such that  $\gamma_{v_1v_2} \neq 0$  and  $v_1, v_2 \in \partial G$ . We note that if we remove a boundary-to-boundary edge, we may be left with a disconnected graph, but that doesn't matter, since we can define medial and dual graphs for disconnected graphs. Similarly all of the results about the minimal boundary value maps were not dependent on the graph being connected.

**Lemma 2.6.4.** Let  $\Gamma$  be an electrical network with boundary-to-boundary edge  $v_1v_2$  and suppose  $\Gamma_0$  is the network resulting from removing this edge. There is an obvious map between the vertices in these networks, which we will call  $\beta : V(G_0) \to V(G)$  which is essentially just the identity. Let  $\phi \in Z(\Gamma)$ . Then  $\phi \in M(\Gamma)$  iff  $(\phi \circ \beta) \in M(\Gamma_0)$ .

*Proof.* Suppose  $\phi \in M(\Gamma) = W(G)^{\perp}$  and let  $u \in W(G_0)$ . Notice that  $\beta$  is a bijection and  $\beta(W(G)) = W(G_0)$ . Hence  $(\phi, u \circ \beta^{-1})_{Z(\Gamma)} = 0$  by assumption. Denote  $u \circ \beta^{-1}$  by  $\widetilde{u}$ . We simply note that since  $\widetilde{u}(v_1) = \widetilde{u}(v_2)$  since  $v_1$  and  $v_2$ 

are boundary vertices. Hence

$$\begin{split} (\phi, \widetilde{u})_{Z(\Gamma)} &= \sum_{V(G) \times V(G)} \gamma_{vv'}(\phi(v) - \phi(v'))(\widetilde{u}(v) - \widetilde{u}(v')) \\ &= \sum_{(V(G) \times V(G)) \setminus \{v_1 v_2, v_2 v_1\}} \gamma_{vv'}(\phi(v) - \phi(v'))(\widetilde{u}(v) - \widetilde{u}(v')) \\ &= \sum_{V(G_0) \times V(G_0)} \gamma_{vv'}((\phi \circ \beta)(v) - (\phi \circ \beta)(v'))(u(v) - u(v')) \\ &= (\phi \circ \beta, u)_{Z(\Gamma_0)} \end{split}$$

and hence  $(\phi \circ \beta, u)_{Z(\Gamma_0)} = 0$ . Hence  $\phi \circ \beta \in Z(\Gamma_0)$ . The other direction follows by reversing the order we presented the above inequalities.

**Lemma 2.6.5.** Given the minimal Dirichlet-to-Neumann map  $\Lambda_M(\Gamma)$  and a boundary-to-boundary edge  $v_1v_2$  with conductance  $\gamma_{v_1v_2}$ , we can find the minimal Dirichlet-to-Neumann map  $\Lambda_M(\Gamma_0)$  for the connected network  $\Gamma_0$  resulting from the removal of the edge  $v_1v_2$ .

*Proof.* By Lemma 2.6.4 we know that the minimal voltage functions on  $\Gamma_0$  are the same as they are on  $\Gamma$ , and hence given valid boundary data, the Dirichlet solutions are the same. If  $\phi$  is a voltage function, by definition, the current leaving a boundary vertex v is given by the formula

$$\Lambda_M(\Gamma)(\phi)(v) = \sum_{v' \sim_G v} \gamma_{vv'}(\phi(v) - \phi(v')).$$

If v is not  $v_1$  or  $v_2$  this is unchanged and hence  $\Lambda_M(\Gamma)(\phi)(v) = \Lambda_M(\Gamma_0)(v)$ . For  $v_1$  and  $v_2$ , we just compute immediately that

$$\begin{split} \Lambda_M(\Gamma)(\phi)(v_1) &= \sum_{v' \sim_G v_1} \gamma_{v_1 v'}(\phi(v_1) - \phi(v')) \\ &= \gamma_{v_1 v_2}(\phi(v_1) - \phi(v_2)) + \sum_{v' \sim_G v_1} \gamma_{v_1 v'}(\phi(v_1) - \phi(v')) \\ &= \Lambda_M(\Gamma_0)(v_1) + \gamma_{v_1 v_2}(\phi(v_1) - \phi(v_2)), \end{split}$$

which gives us a formula for  $\Lambda_M(\Gamma_0)(v_1)$ . Using an identical argument we can find a nearly identical formula for  $\Lambda_M(\Gamma_0)(v_2)$ .

**Lemma 2.6.6.** Given a critical half planar electrical network  $\Gamma$  with a boundaryto-boundary edge  $v_1v_2$ , the map  $\Lambda_M(\Gamma)$  uniquely determines  $\gamma_{v_1v_2}$ .

*Proof.* The proof is essentially the same as in the case of boundary spikes because now we are recovering a boundary spike of the dual graph. We leave the details to the reader, but essentially we find a geodesic g which crosses the boundary spike in the dual graph and such that the boundary vertex of this

boundary spike is in B(g). We find boundary values which will guarantee that the interior vertex of this boundary spike will have covoltage zero but such that the boundary vertex will have covoltage 1. We leave the details to the reader since they are identical to before.

## 2.6.3 Recovering the Entire Graph

**Definition 2.6.7.** We will define a **triangular geodesic region** to be a bounded subset of the medial graph whose boundary consists of two exactly geodesic segments and one connected subset of  $\mathbb{R}$ . We will define a **geodesic triangle** to be a triangular geodesic region R such that there are no geodesics which cross the boundary of R.

**Lemma 2.6.8.** Let g be a geodesic in a critical half planar medial graph and suppose that g intersects at least one other geodesic. Then there is a geodesic triangle (possibly with other geodesics inside of it) in B(g).

*Proof.* The proof is essentially the same as in [1]. Let  $x_{\ell}^g$  and  $x_r^g$  be the left and right endpoints of g. Let g' be the geodesic which crosses g closest to  $x_{\ell}$ . One of the endpoints of g' must be in B(g) since the graph is critical. This produces a triangular region t. Let g'' be the geodesic which crosses g' closest to this region. Notice that g'' cannot intersect g' again and g'' cannot intersect g at any point in  $\bar{t}$  since we assumed that g' was the geodesic which was closest to x. This yields a geodesic triangle  $t' \subseteq t \subseteq B(g)$ . We summarize in Figure 2.20.



Figure 2.20: A decreasing sequence of triangular geodesic regions in B(q).

We repeat this process to get a descending sequence of triangular geodesic regions which must eventually terminate with a geodesic triangle since B(g) is a finite subset of the medial graph. We note that we don't claim the resulting geodesic triangle is empty.

**Lemma 2.6.9.** Removal of a boundary-to-boundary edge, a boundary spike, or removing a finite connected component all preserve criticality.

*Proof.* All operations obviously perserve half-planarity. The first two operations preserve criticality since the change in the medial corresponds to just merging two cells as shown in Figure 2.21.



Figure 2.21: The result on the medial graph from a boundary spike contraction. The result of a boundary-to-boundary edge deletion is the same since the two processes are dual to each other.

And hence we will never have two geodesics if they didn't cross before removing a boundary spike or boundary to boundary edge. If there is a finite connected component of the graph, then one there must be a dual cell which has multiple boundary components on  $\mathbb{R}$  and hence deleting the finite subset will just correspond to deleting some geodesics, as shown in the below picture. We leave the details to the reader, but the claim is clear.

**Lemma 2.6.10.** Let G be a critical half planar graph and let e be any edge in G. Then e can be removed by sequentially removing boundary spikes, boundary-to-boundary edges or deleting finite connected components.

**Proof.** Let g be a geodesic which crosses the edge e. we will sequentially remove every edge with a vertex in B(g) (corresponding to removing a cell in the medial graph). So we proceed by induction, showing that we can always remove cells in B(g) or merge cells in B(g) without altering any of the rest of the medial graph. To do this, we note that since e crosses g, there must by another geodesic which crosses e and hence crosses g. Thus we can apply Lemma 2.6.8 to find a geodesic triangle t in B(g). If t is empty, then we can simply perform remove a boundary spike or boundary to boundary edge, which results in just merging two cells as shown in Figure 2.21.

As shown in Lemma 2.6.9 this preserves criticality. If t is nonempty then since there are no geodesics which cross t, we know that the vertices of G in t cannot be connected to any of the vertices in  $M \setminus t$ , and hence the vertices in t correspond to a connected component so we can remove them. As shown in Lemma 2.6.9, we know that this preserves criticality and deleting the primal vertices in t correspond to removing the geodesics inside of t which don't cross t. Eventually we must get that no geodesics cross g, which implies that we have removed e.

**Theorem 2.6.11.** Given a critical half planar network graph G and the minimal Dirichlet-to-Neumann map  $\Lambda_M$  for some electrical network  $\Gamma$  with graph G, we can recover all of the conductivities of  $\Gamma$ .

**Proof.** Let e be an edge in G. By Lemma 2.6.10 we can remove e from the graph by sequentially removing boundary spikes, boundary-to-boundary edges and finite connected components. Given a boundary spike, or a boundary-to-boundary edge we can recover the conductivity by Lemmas 2.6.3 and 2.6.6 we can recover the conductivity along it. By Lemmas 2.6.2 and 2.6.5 we can find the Dirichlet to Neumann maps of the resulting graphs. Given a finite connected component of G, the Dirichlet-to-Neumann map for the finite component is just the restriction of the minimal Dirichlet-to-Neumann map for the infinite graph, and since the finite component must obvious be critical circular planar we can recover all of the conductivities of the edges in the finite component. The minimal Dirichlet-to-Neumann map for the other components will just be the restriction of the Dirichlet-to-Neumann map restricted to the complement of the finite component. Repeating this process as per Lemma 2.6.10, we will recover the entire graph.

We note that there there is an obvious analogue for Lemmas 2.6.2, and 2.6.5 for the minimal Neumann-to-Dirichlet map  $H_M$ . Similarly, the results from Lemmas 2.6.3 and 2.6.6 hold for  $\Lambda_M$  replaced with  $H_M$  since the proof basically carries over without change. Thus we have the proposition, the details are left the reader, but are essentially just as above:

**Proposition 2.6.12.** Given a critical half planar network graph G and the minimal Neumann-to-Dirichlet map  $H_M$  for some electrical network  $\Gamma$  with graph G, we can recover all of the conductivities of  $\Gamma$ .

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